

# Particle Creation in presence of A Warped Extra Dimension

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# Plan of Talk

- Introduction
- Quantum field coupled to a background metric giving rise to particle creation and constraints on the field modes
- Production of particles in presence of different kinds of extra dimension and their generic features
- Discussions

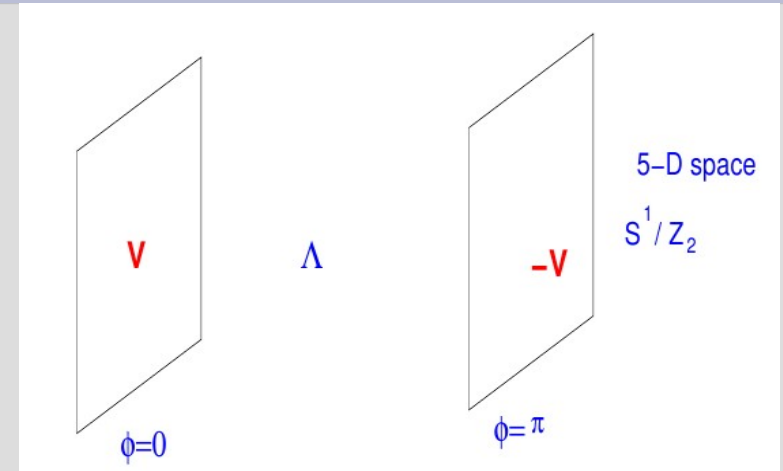
# Introduction

## The Braneworld scenario

*L. Randall & R. Sundrum, 1999*

$$ds^2 = e^{2f(\sigma)} \eta_{\mu\nu} dx^\mu dx^\nu + d\sigma^2$$

- Our world is a four dimensional hypersurface (3-brane) embedded in higher dimensional universe.
- The warping factor.



RS II scenario

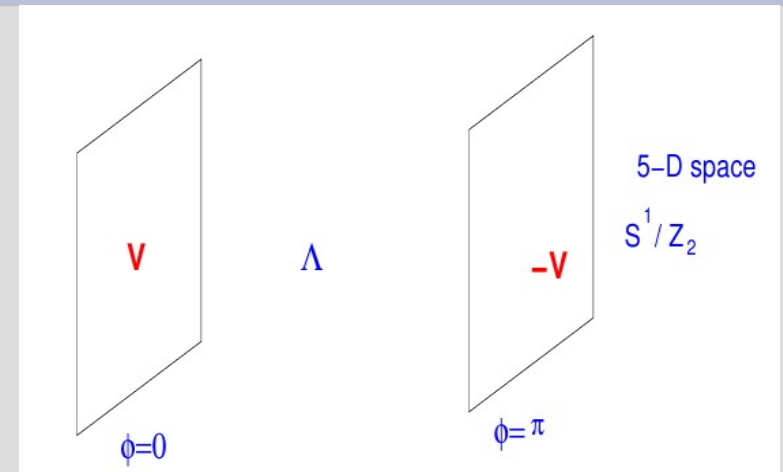
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These warped Braneworld models propose:

achievable experiments in search for signature of extra dimensions,  
a solution to long standing hierarchy problem in high energy physics,  
a dynamic way of compactification by localisation of fields on Brane  
and many more interesting features to be explored.

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- Quantum fields in warped spacetimes.
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- Quantum fields in warped spacetimes
- Particle creation can be viewed as a quantum consequence of extra dimension.

The questions arise are:

- How do the extra dimension effect particle production ?
- What is the role of warping factor, cosmological evolution and time dependence of extra dimension?
- Is there any new feature ?

# Quantum field coupled to a spacetime with an extra dimension

## A formalism

A scalar field  $\psi(\eta, \mathbf{x}, \sigma)$  Conformally coupled with

The metric  $ds^2 = e^{2f(\sigma)} a^2(\eta) [-d\eta^2 + dx^2 + dy^2 + dz^2] + \phi^2(\eta) d\sigma^2$

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gives rise to the following Klein-Gordon equation

$$-\frac{e^{-2f(\sigma)}}{a^2} \frac{\partial^2 \psi}{\partial \eta^2} + \frac{e^{-2f(\sigma)}}{a^2} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + \frac{1}{\phi^2} \frac{\partial^2 \psi}{\partial \sigma^2} - \frac{e^{-2f(\sigma)}}{a^2} \left( \frac{2\dot{a}}{a} + \frac{\dot{\phi}}{\phi} \right) \frac{\partial \psi}{\partial \eta} + \frac{4f'}{\phi^2} \frac{\partial \psi}{\partial \sigma} - (m^2 + \xi R) \psi = 0$$



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Assuming separability  $\longrightarrow \psi(\eta, \mathbf{x}, \sigma) = \frac{1}{e^{f(\sigma)} a \phi^{\frac{1}{2}}} \chi_l(\eta) F(\mathbf{x}) G(\sigma)$

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$$\ddot{\chi}_l(\eta) + \left[ \left( \mathbf{k}^2 + \frac{a^2}{\phi^2} k_\sigma^2 \right) + \frac{\ddot{a}}{8a} - \frac{\ddot{\phi}}{8\phi} + \frac{\dot{\phi}^2}{4\phi^2} - \frac{\dot{a}\dot{\phi}}{4a\phi} \right] \chi_l(\eta) = 0 \quad m=0 \text{ and } \xi = \frac{3}{16}$$

$$\frac{1}{F(\mathbf{x})} \left\{ \frac{d^2 F(\mathbf{x})}{dx^2} + \frac{d^2 F(\mathbf{x})}{dy^2} + \frac{d^2 F(\mathbf{x})}{dz^2} \right\} = -\mathbf{k}^2$$

$$e^{2f} \left\{ \frac{G''(\sigma)}{G(\sigma)} + 2f' \frac{G'(\sigma)}{G(\sigma)} + \left( \frac{f''}{2} + \frac{3f'^2}{4} \right) \right\} = -k_\sigma^2$$

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In general  $\rightarrow$  
$$\ddot{\chi}_l(\eta) + [\Omega_l^2 + Q] \chi_l(\eta) = 0$$

wronskian - 
$$\dot{\chi}_l^* \chi_l - \dot{\chi}_l \chi_l^* = i$$

$$\Omega_l^2 = \left( \mathbf{k}^2 + \frac{a^2}{\phi^2} k_\sigma^2 \right)$$

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WKB solutions  $\rightarrow$

$$\chi_l = \frac{\alpha_l}{\sqrt{2\Omega_l}} e^{-i \int^\eta \Omega_l d\eta} + \frac{\beta_l}{\sqrt{2\Omega_l}} e^{i \int^\eta \Omega_l d\eta}$$

further restriction -

$$\dot{\chi}_l = -i\Omega_l \left[ \frac{\alpha_l}{\sqrt{2\Omega_l}} e^{-i \int^\eta \Omega_l d\eta} - \frac{\beta_l}{\sqrt{2\Omega_l}} e^{i \int^\eta \Omega_l d\eta} \right]$$

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$$\dot{\alpha}_l = \frac{1}{2} \left( \frac{\dot{\Omega}_l}{\Omega_l} - i \frac{Q}{\Omega_l} \right) \beta_l e^{+2i \int \Omega_l d\eta} - i \frac{Q}{2\Omega_l} \alpha_l$$

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## Zel'dovich & Starobinsky

$$\frac{ds_l}{d\eta} = \frac{\dot{\Omega}_l}{2\Omega_l} v_l + \frac{Q}{2\Omega_l} r_l,$$

$$\frac{dv_l}{d\eta} = \frac{\dot{\Omega}_l}{\Omega_l} (1 + 2s_l) - \left[ \frac{Q}{\Omega_l} + 2\Omega_l \right] r_l$$

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$$s_l(\eta_0) = r_l(\eta_0) = v_l(\eta_0) = 0$$

$$s_l \equiv |\beta_l|^2$$

$$v_l \equiv 2 \operatorname{Re} \left[ \alpha_l \beta_l^* e^{-2i \int^\eta \Omega_l d\eta} \right]$$

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$$N = \frac{1}{(2\pi a)^3} \int d^3l |\beta_l|^2$$

Total particle number density

$$\rho = \frac{1}{(2\pi a)^3 a} \int d^3l l |\beta_l|^2$$

Total energy density

# Allowed values of $k_\sigma$

Eigenvalue equation  $\rightarrow$

$$e^{2f} \left\{ \frac{G''(\sigma)}{G(\sigma)} + 2f' \frac{G'(\sigma)}{G(\sigma)} + \left( \frac{f''}{2} + \frac{3f'^2}{4} \right) \right\} = -k_\sigma^2$$

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We get  $G_1''(\sigma) + \left[ b \delta(\sigma) - \frac{b^2}{4} + k_\sigma^2 e^{2b|\sigma|} \right] G_1(\sigma) = 0$  for  $f(\sigma) = -b|\sigma|$

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For a two brane model:

$$\tan \left[ \theta + \tan^{-1} \left( -\frac{2}{3} \theta e^{-br_c \pi} \right) \right] = -\frac{2}{3} \theta \quad \text{where} \quad k_\sigma = \theta b e^{-br_c \pi}$$

$$\theta = 0, \pm 2.17463, \pm 5.00365, \pm 8.03846, \pm 11.1295, \pm 14.2421, \pm 17.3649 \text{ and so on}$$

In case of an infinite extra dimension there are no discrete values of the modes and we have a continuum.

# Particle creation in a four dimensional universe

4d metric  $ds^2 = a(\eta)^2(-d\eta^2 + dx^2 + dy^2 + dz^2)$  with  $a^2(\eta) = b_1^2 + b_2^2\eta^2$



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K-G eq.  $\ddot{\chi}_l(\eta) + \Omega_l^2(\eta)\chi_l(\eta) = 0$

where  $\Omega_l^2(\eta) = k^2 + (b_1^2 + b_2^2\eta^2)m^2$

$$\ddot{\chi}_l(\eta) + [\Omega_l^2 + Q]\chi_l(\eta) = 0$$

$Q(\eta)$  is the signature of breaking of conformal invariance.

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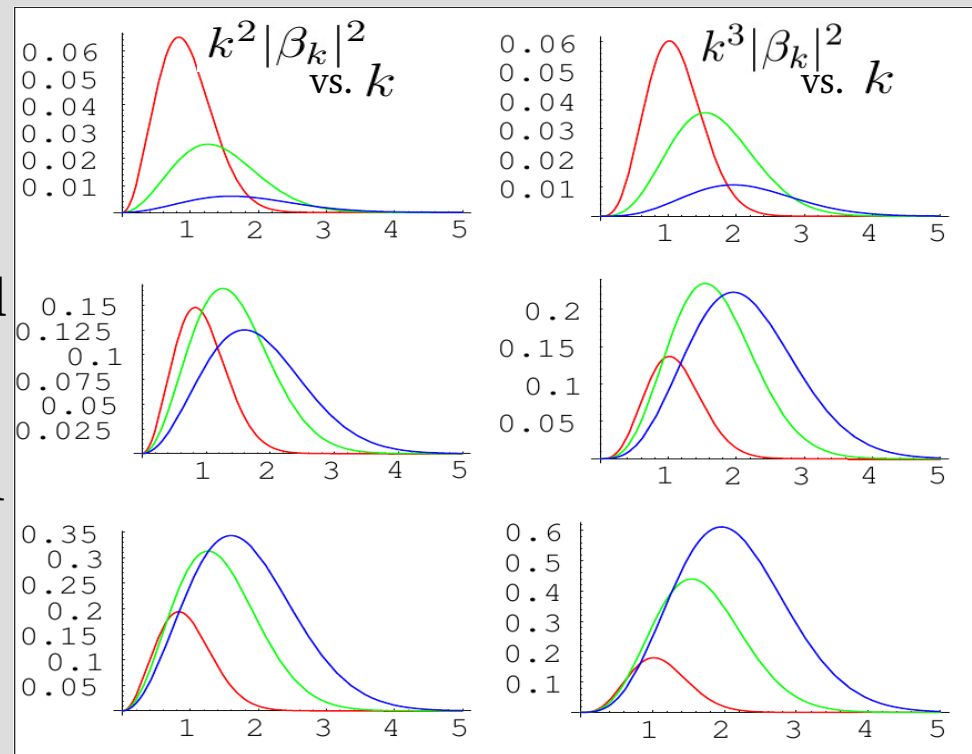
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$$|\beta_k|^2 = \exp \left[ -\pi \left( \frac{k^2}{mb_2} + \frac{mb_1^2}{b_2} \right) \right]$$

$k^2|\beta_k|^2 \rightarrow$  Integrand to compute total particle number density

$k^3|\beta_k|^2 \rightarrow$  Integrand to compute total energy density



Decreasing  $b_1 \downarrow$

# Particle creation in a universe with a static extra dimension

$$ds^2 = e^{2f(\sigma)} a^2(\eta) [-d\eta^2 + dx^2 + dy^2 + dz^2] + \phi^2(\eta) d\sigma^2 \quad a^2(\eta) = b_1^2 + b_2^2 \eta^2 \quad \phi(\eta) = b_1$$

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$$\ddot{\chi}_l(\eta) + \left[ \left( k^2 + \frac{b_1^2 + b_2^2 \eta^2}{b_1^2} k_\sigma^2 \right) + \frac{b_1^2 b_2^2}{8(b_1^2 + b_2^2 \eta^2)} \right] \chi_l(\eta) = 0$$

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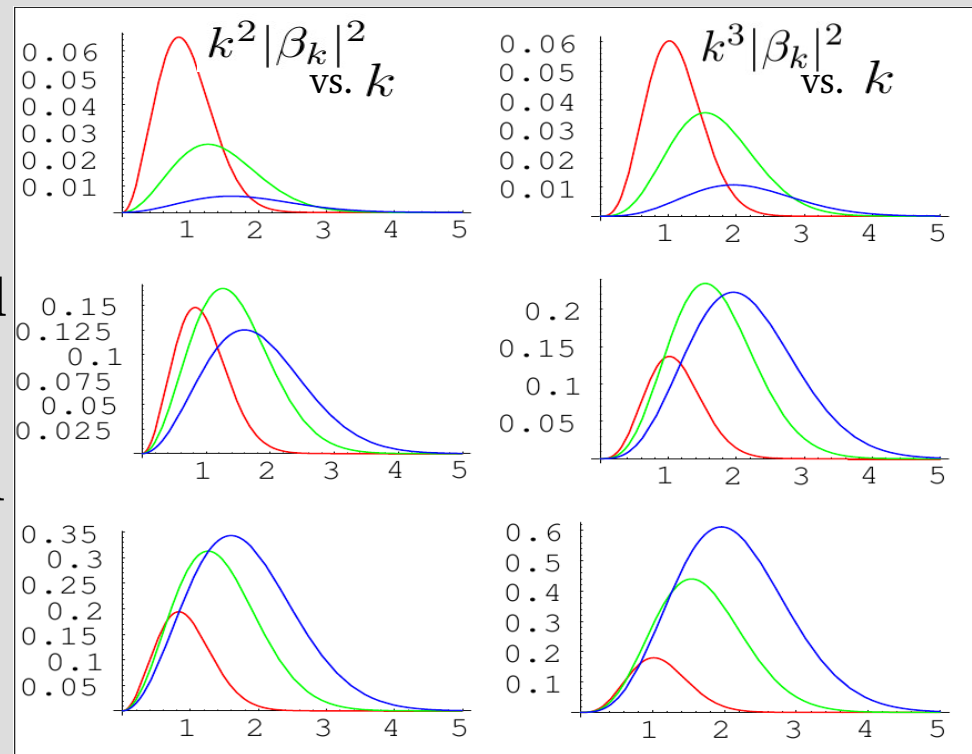
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$$\ddot{\chi}_l(\eta) + \left[ \left( k^2 + \frac{b_1^2 + b_2^2 \eta^2}{b_1^2} k_\sigma^2 \right) + \frac{b_1^2 b_2^2}{8(b_1^2 + b_2^2 \eta^2)} \right] \chi_l(\eta) = 0 \quad Q(\eta)=0 \Rightarrow m \equiv \frac{k_\sigma}{b_1}$$

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Numerical evaluation

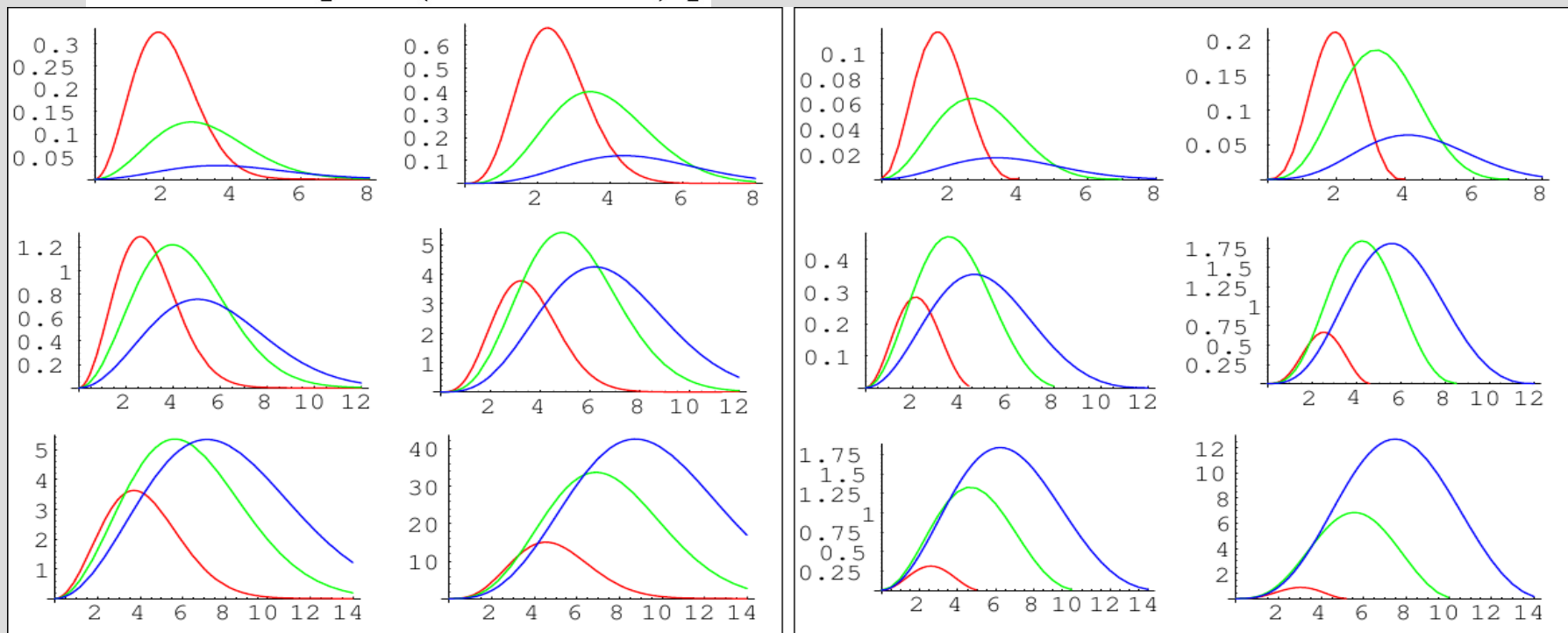
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$$k_\sigma = 2.17463, 5.00365, 8.03846$$

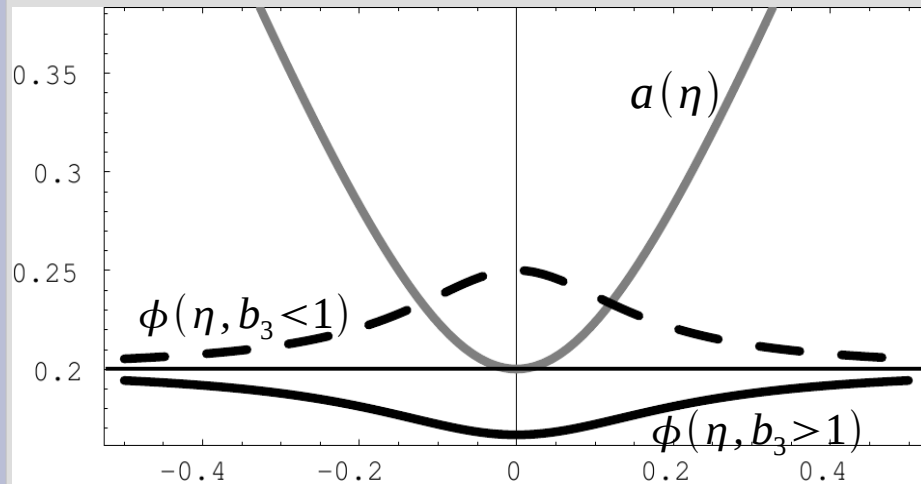


# Particle creation in a universe with a time-dependent extra dimension

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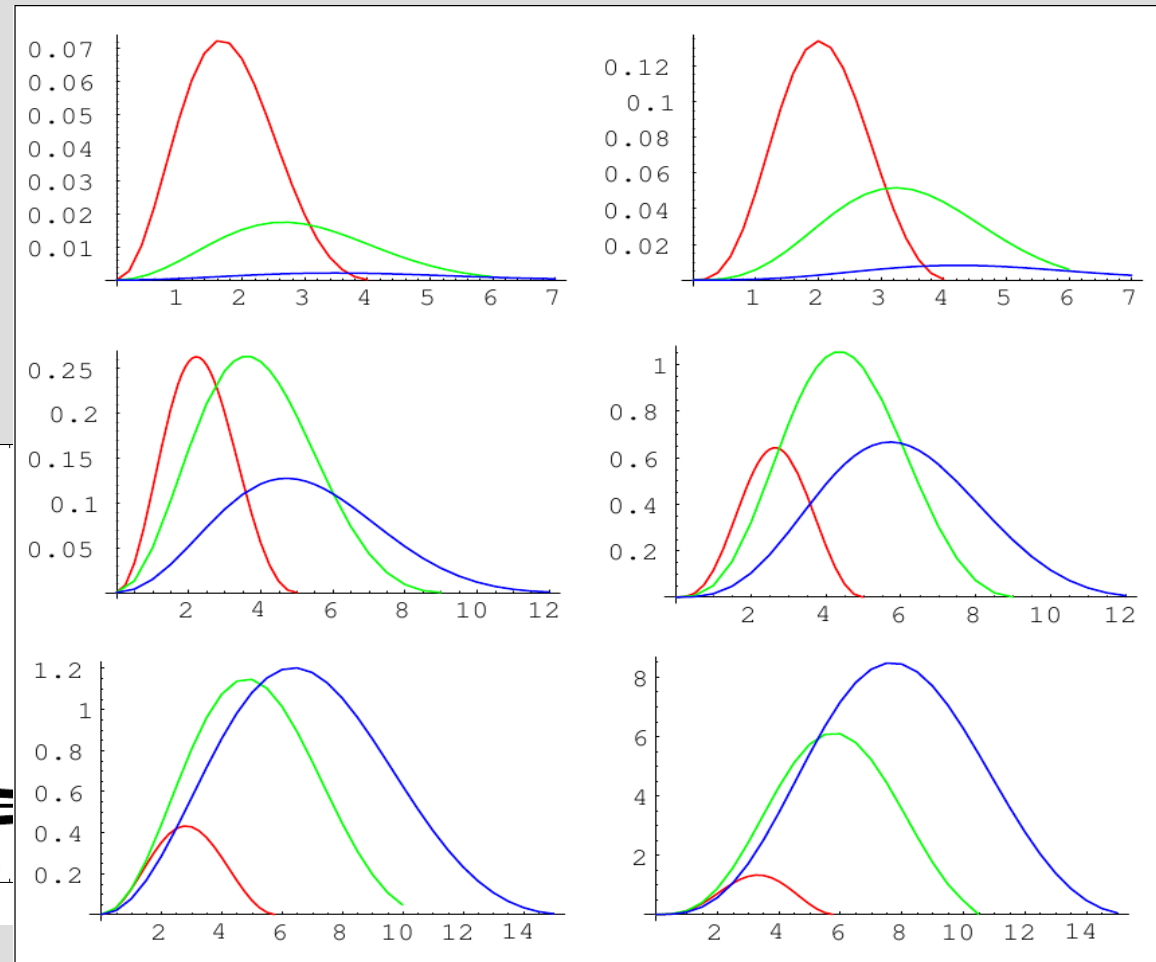
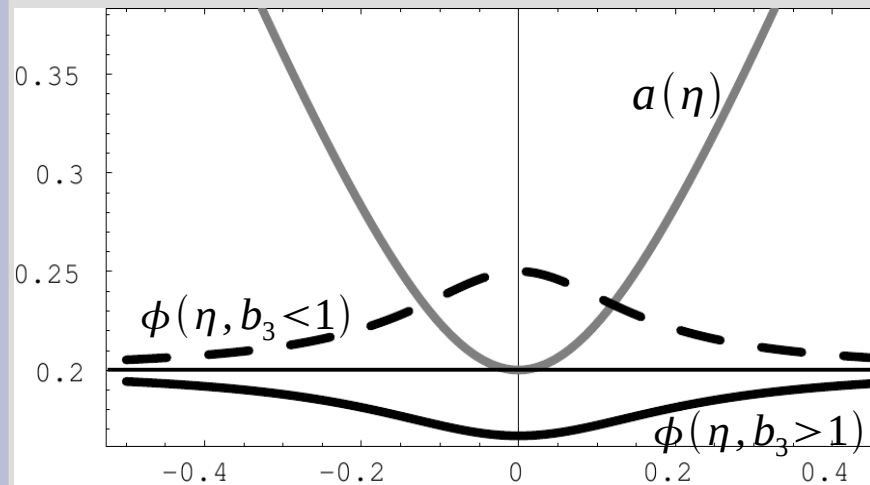


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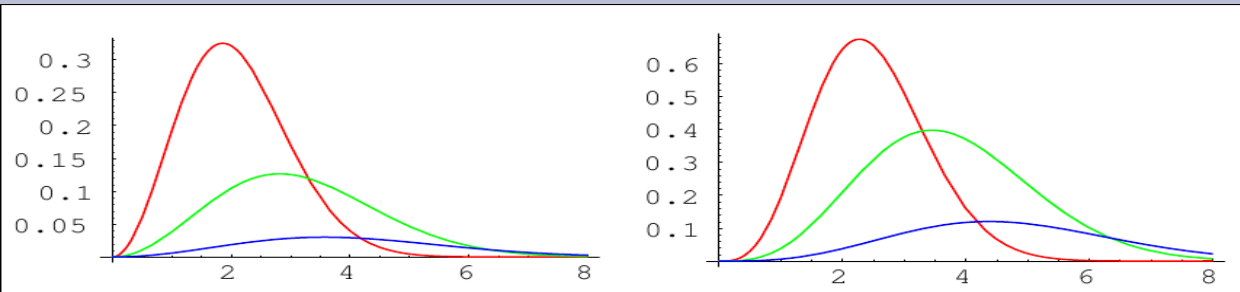


$\phi(\eta, b_3 > 1)$

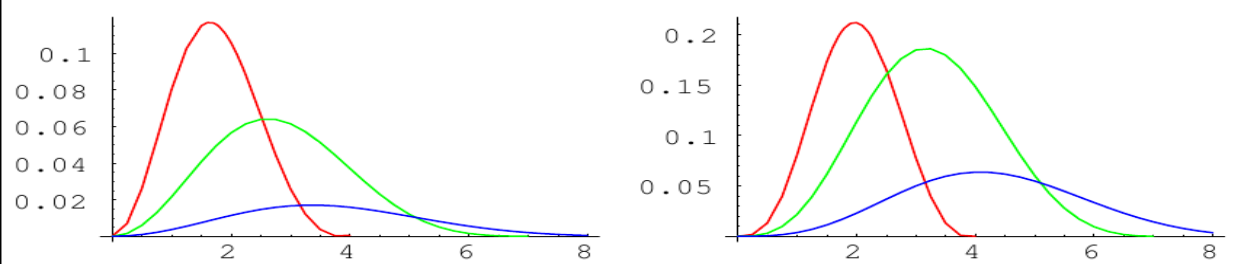
# Comparison

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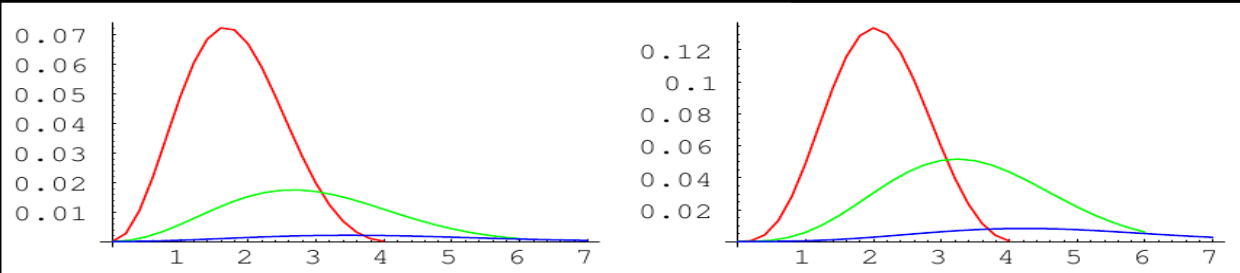
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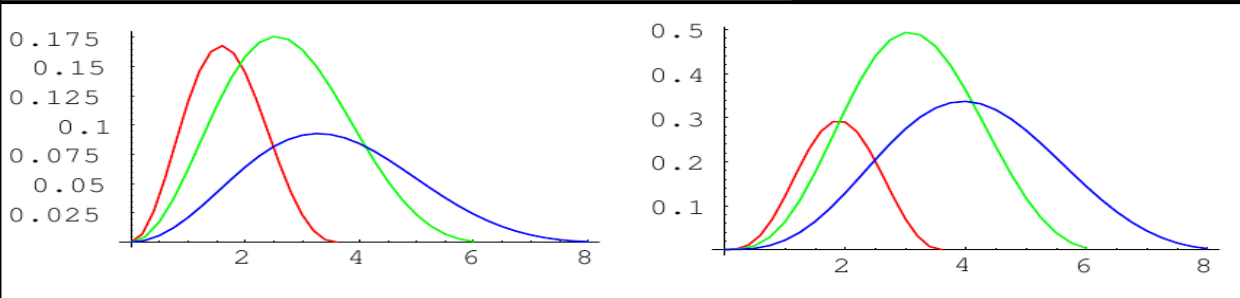
Static extra dimension  
Analytic approximation  
 $Q(\eta)=0$



Static extra dimension  
Numerical evaluation  
 $Q(\eta) \neq 0$



$$\phi^2(\eta) = \frac{b_1^2 + b_2^2 \eta^2}{b_3^2 + b_2^2 \eta^2 / b_1^2} \quad \text{and} \quad b_3 > 1$$



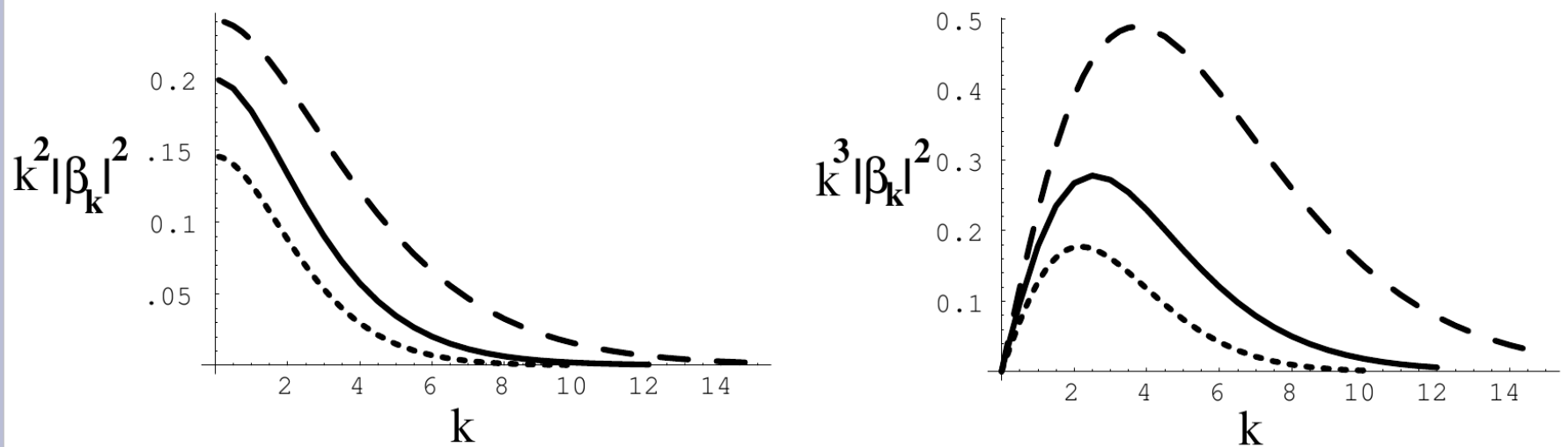
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Unlike in case of a radiative universe, conformal invariance is broken even for the massless (4d) modes.

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Continuous curve: Static extra dimension

Dashed curve: Dynamic extra dimension with  $b_3 < 1$

Dotted curve: Dynamic extra dimension with  $b_3 > 1$

# Summary & comments

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- Effects of  $Q(\eta)$  :

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- One can try to derive such braneworld type solutions of Einsteins equation that we have assumed as our model.
- This same study can be done for different cosmological models.
- One can carry out the same study for different kinds of particles.

*Thank You*