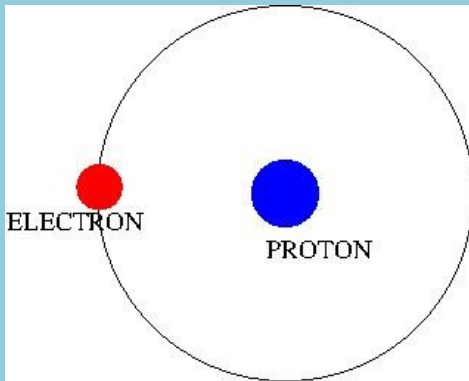


# 21-cm Physics and Cosmology

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# Hydrogen atom



Spin of proton and electron  
Is ignored

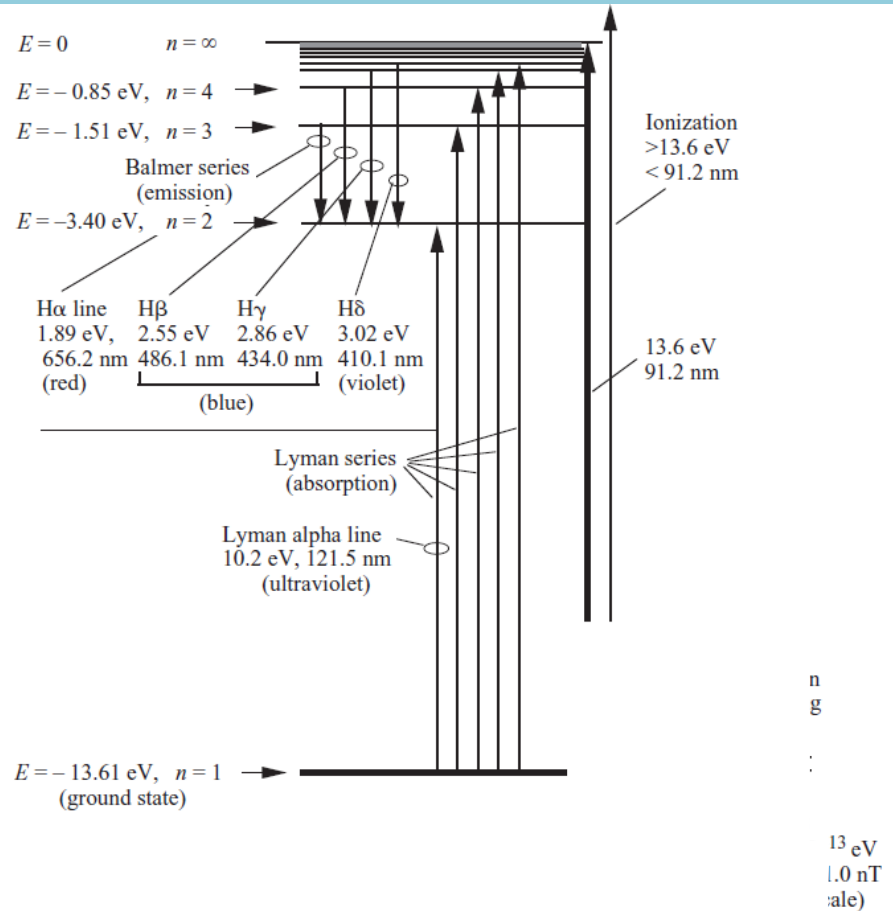
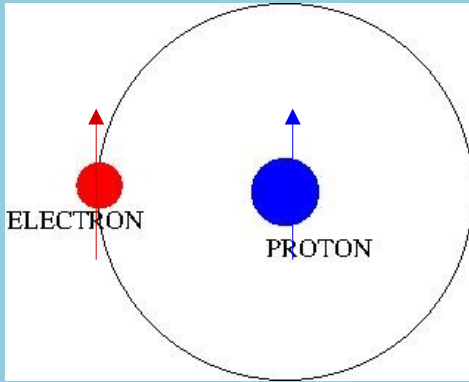


Fig. 10.3: Energy levels of the hydrogen atom. Downward Balmer and upward Lyman transitions are shown. The minuscule hyperfine splitting of the ground state is shown with greatly exaggerated spacing together with the even finer Zeeman splitting of the upper hyperfine level. The energy difference between the upper and lower Zeeman states is given for a 1-nT line-of-sight component of magnetic field in a hydrogen cloud. [Adapted from H. Bradt, *Astronomy Methods*, Cambridge, 2004, Fig. 101, with permission]

# Ground State n=1



Has zero orbital angular momentum

Total angular momentum is sum of electron and proton spins

$$F = S + I.$$

Electro

$$S = 1/2$$

$$S \cdot S \Rightarrow S(S + 1)\hbar^2.$$

$$S_z \Rightarrow m_S \hbar; m_S = +1/2, -1/2.$$

Proton

$$I = 1/2$$

$$I \cdot I \Rightarrow I(I + 1)\hbar^2$$

$$I_z \Rightarrow m_I \hbar; m_I = +1/2, -1/2.$$

Triplet

UU

$$F = 1$$

$$m_F = +1$$

$$\frac{UD + DU}{\sqrt{2}}$$

$$F = 1$$

$$m_F = 0$$

DD

$$F = 1$$

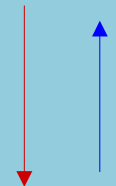
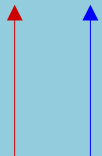
$$m_F = 1$$

$$\frac{UD - DU}{\sqrt{2}}$$

$$F = 0$$

$$m_F = 0.$$

Singlet



# Magnetic Moments

Electro

$$\mu_e = g_e \frac{\mu_B}{\hbar} S$$

Bohr Magneton

$$\mu_B \equiv \frac{e \hbar}{2 m_e} = 9.274 01 \times 10^{-24} \text{ J/T}$$

$$g_e = -2.00232$$



Proto

$$\mu_p = g_p \frac{\mu_N}{\hbar} I$$

Nuclear Magneton

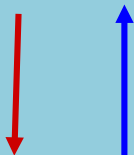
$$\mu_N \equiv \frac{e \hbar}{2 m_p} = 5.050 783 \times 10^{-27} \text{ J/T}$$

$$g_p = +5.586$$

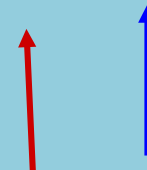


Proton magnetic moment is  
~650 times smaller than electron

Triplet



Singlet



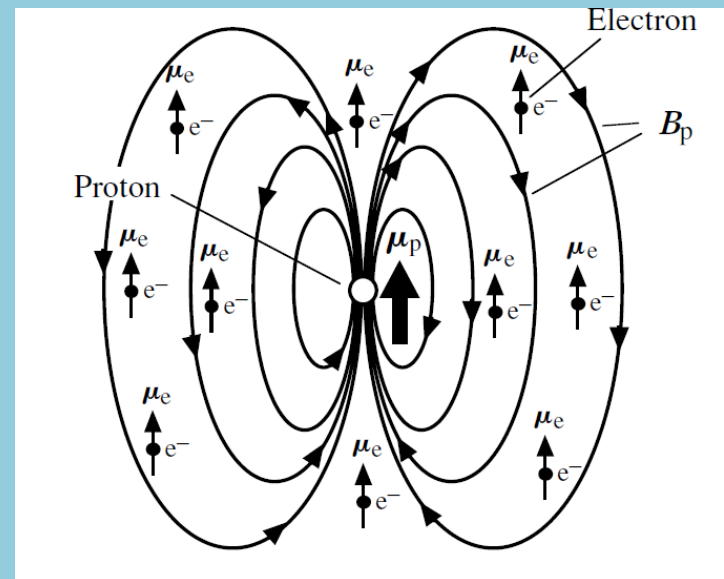
# Interaction Energy

$$E_{\text{pot}} = -\boldsymbol{\mu} \cdot \mathbf{B}$$

1. Proton's Magnetic Field and Electron's Orbital Magnetic Moment – Not there in ground state.

2. Proton's Magnetic Field and Electron's Spin Magnetic Moment – Zero when averaged over Electron's Spatial Distribution

$$\psi(r) = \frac{1}{\sqrt{4\pi}} \frac{2}{a_0^{3/2}} \exp(-r/a_0).$$



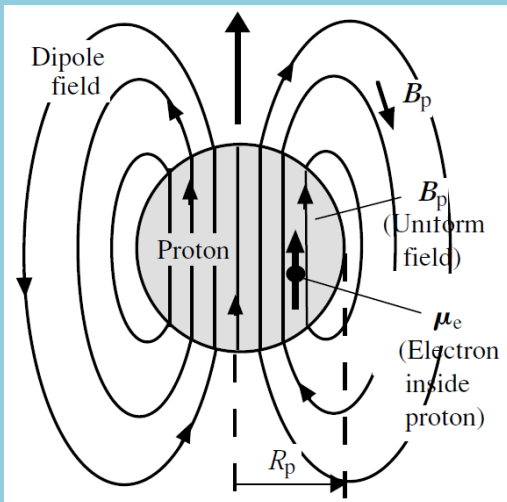
Bohr Radius

$$a_0 = 5.291772 \times 10^{-11} \text{ m,}$$

# Contact Term

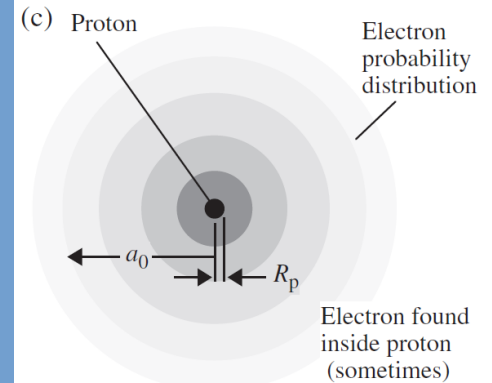
$$\mathbf{B}_p = \frac{\mu_0}{2\pi R_p^3} \boldsymbol{\mu}_p$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ T m A}^{-1}$$



$$E_{\text{pot}} = -\boldsymbol{\mu} \cdot \mathbf{B}$$

$$= -\frac{4}{3} \frac{R_p^3}{a_0^3} \boldsymbol{\mu}_e \cdot \mathbf{B}_p$$



$$-\frac{4}{3} \frac{\mu_0}{2\pi} \frac{g_e g_p \mu_B \mu_N}{a_0^3} \frac{\mathbf{S} \cdot \mathbf{I}}{\hbar^2} \Rightarrow E_{\text{pot}}$$

$$\mathbf{S} \cdot \mathbf{I} \Rightarrow \frac{1}{2} \hbar^2 [F(F+1) - S(S+1) - I(I+1)]$$

$$\mathbf{S} \cdot \mathbf{I} \Rightarrow \frac{1}{2} \hbar^2 \left[ 0 - \frac{3}{4} - \frac{3}{4} \right] = -\frac{3}{4} \hbar^2$$

$$F = 0$$

$$\mathbf{S} \cdot \mathbf{I} \Rightarrow \frac{1}{2} \hbar^2 \left[ 2 - \frac{3}{4} - \frac{3}{4} \right] = -\frac{1}{4} \hbar^2$$

$$F = 1$$

# Hyperfine Splitting

$$\Delta E_{\text{pot}} = E_{\text{pot}, F=1} - E_{\text{pot}, F=0} = 9.42762 \times 10^{-25} \left[ \frac{1}{4} - \left( -\frac{3}{4} \right) \right]$$

1422.8 MHz

More precise QED calculation

$$\nu = 1\,420\,405\,752 \text{ Hz}$$

1420.4 MHz.

$\lambda = 21.11 \text{ cm}$

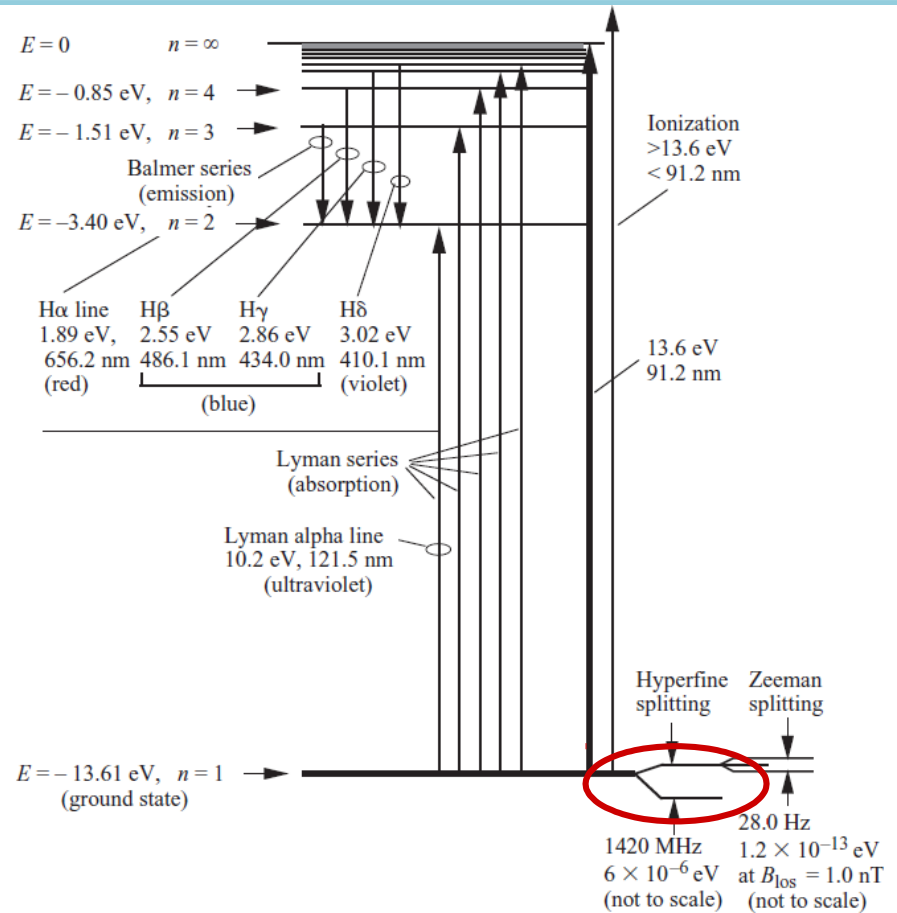
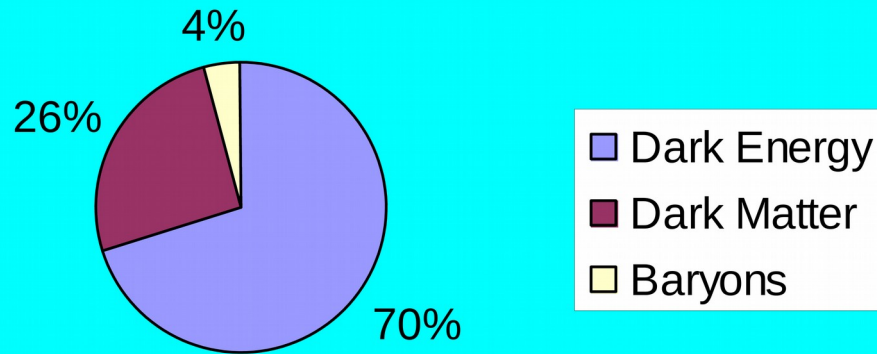


Fig. 10.3: Energy levels of the hydrogen atom. Downward Balmer and upward Lyman transitions are shown. The minuscule hyperfine splitting of the ground state is shown with greatly exaggerated spacing together with the even finer Zeeman splitting of the upper hyperfine level. The energy difference between the upper and lower Zeeman states is given for a 1-nT line-of-sight component of magnetic field in a hydrogen cloud. [Adapted from H. Bradt, *Astronomy Methods*, Cambridge, 2004, Fig. 101, with permission]

# Constituents of the Universe

Present Constituents of Universe



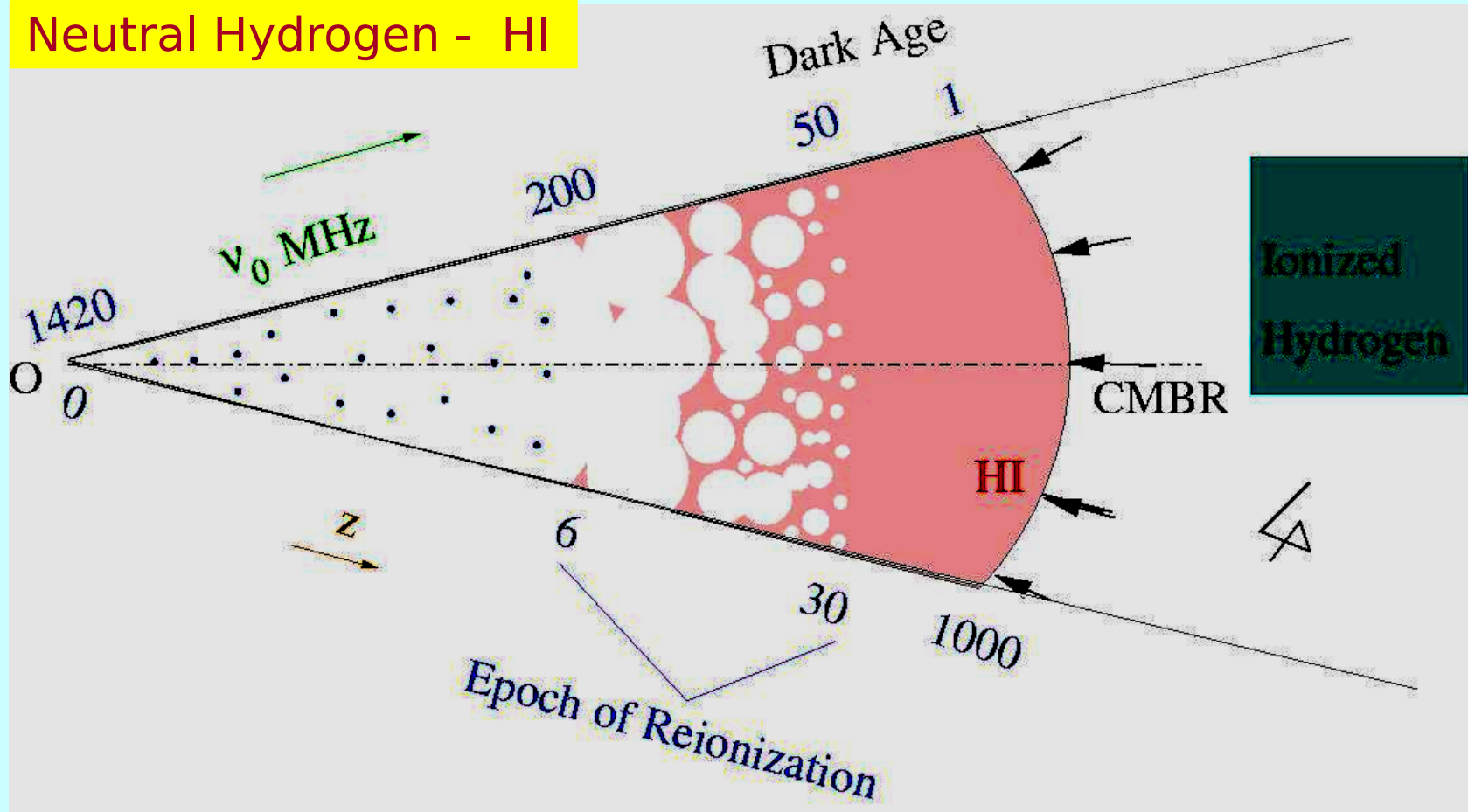
## Baryons

75 %  
Hydrogen

25% Helium

# HI Evolution

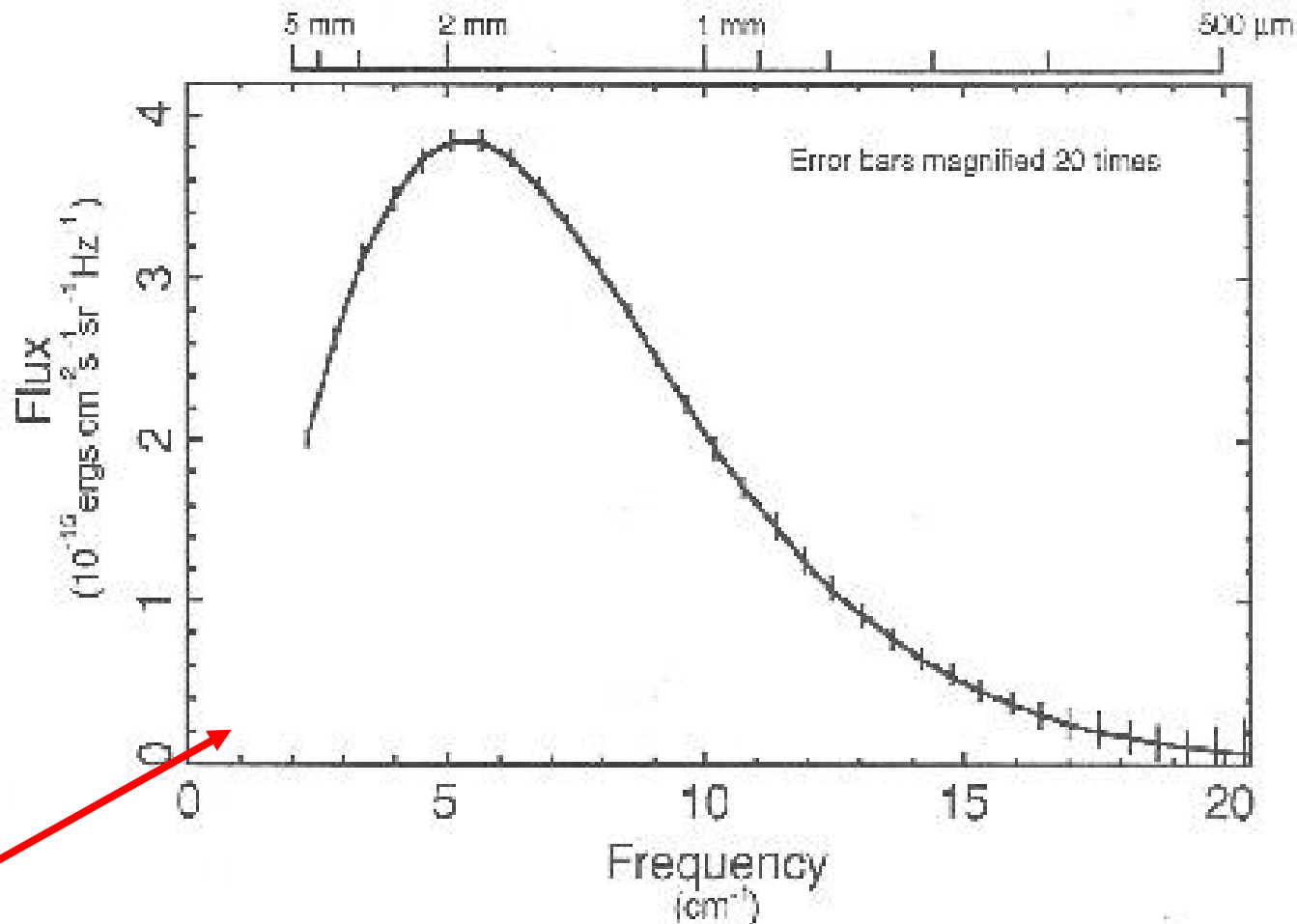
## Neutral Hydrogen - HI



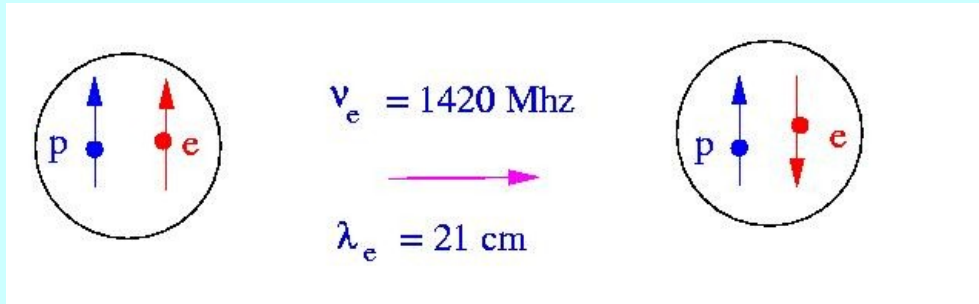
$$\nu_o = 1420 \text{ Mhz} / (1+z)$$

$$\lambda_o = 21 \text{ cm} (1+z)$$

# Cosmic Microwave Background Radiation



# Spin Temperature

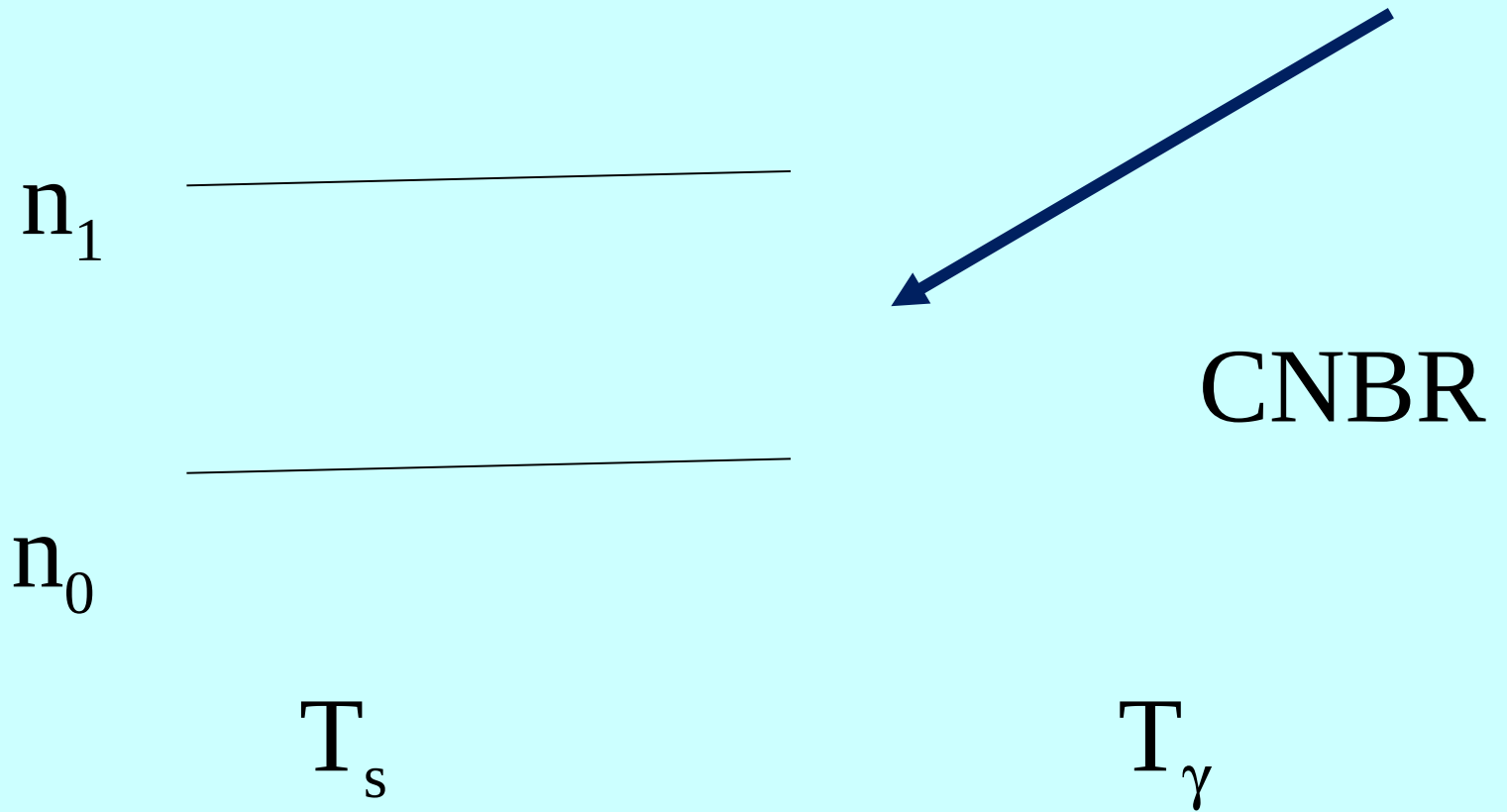


$$\frac{n_1}{n_0} = \frac{g_1}{g_0} e^{-T_\star/T_s}$$

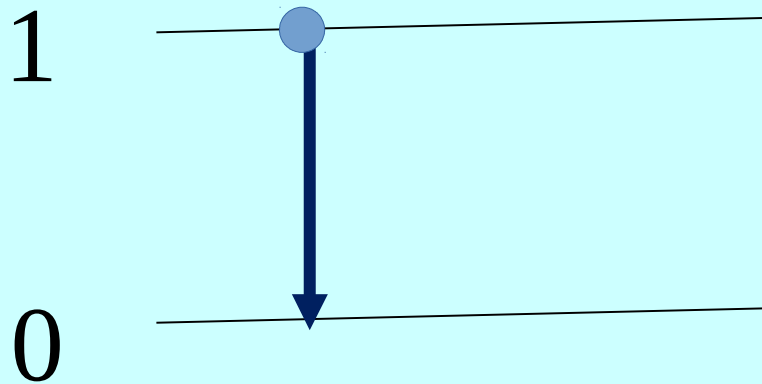
3

$$T_\star = h_p \nu_e / k_B = 0.068 \text{ K}$$

# 21-signal



# Spontaneous Emission



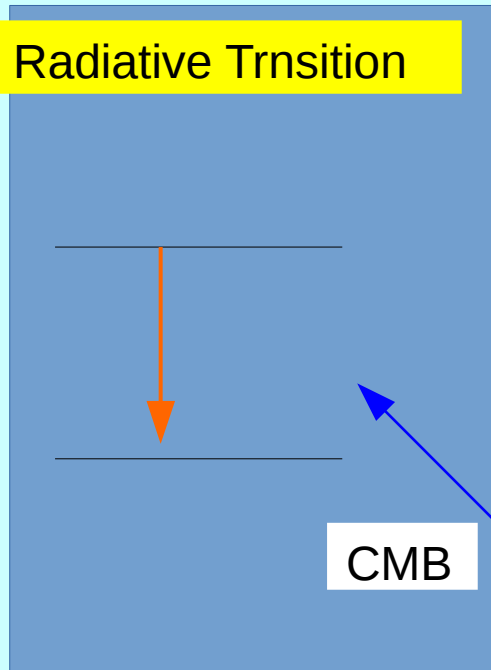
Einstein Coefficient  
Transition probability per unit time

$$A_{10} = 2.85 \times 10^{-15} \text{ s}^{-1}$$

Lifetime ~10 million yrs

# Excitation - De-excitation

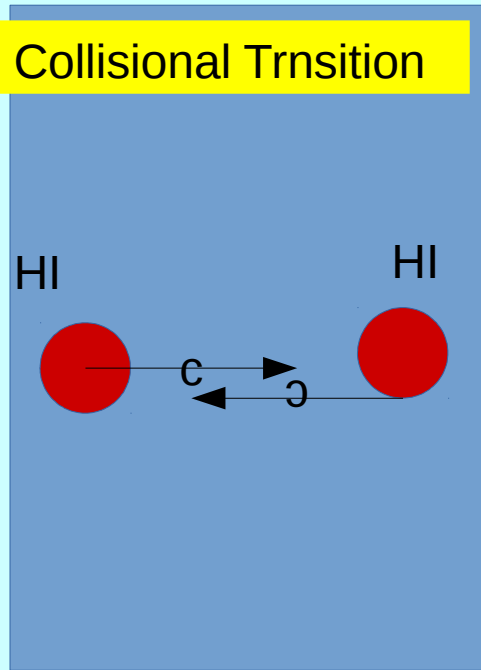
## Radiative Transition



CMB Temperature

$$T_{\gamma}$$

## Collisional Transition

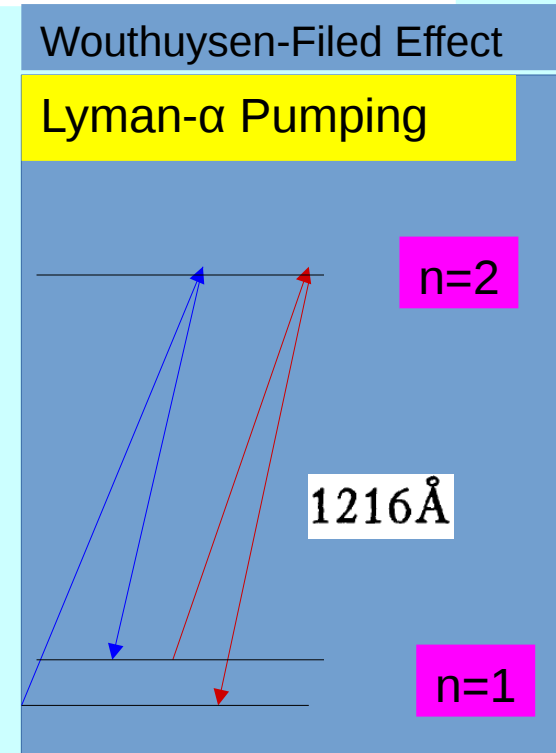


Gas Kinetic Temperature

$$T_K$$

## Wouthuysen-Field Effect

### Lyman- $\alpha$ Pumping



Effective Color Temperature  
Of Lyman- $\alpha$  Radiation

$$\tilde{T}_c \approx T_K$$

# Rate Equation

Excitation rate

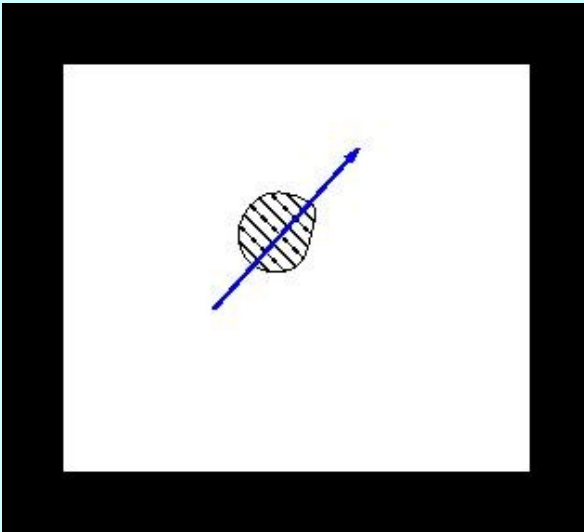
De-excitation rate

$$\dot{n}_0 = -n_0 [ P_{01}^\gamma + P_{01}^c + P_{01}^\alpha ] + n_1 [ P_{10}^\gamma + P_{10}^c + P_{10}^\alpha ]$$

Radiative Collision Lyman- $\alpha$



# Detailed Balance



Consider a situation where radiative processes dominate – in equilibrium

$$T_s = T_\gamma \quad \dot{n}_0 = 0$$

$$n_0 P_{01}^\gamma = n_1 P_{10}^\gamma$$

$$(T_*/T_\gamma) \ll 1$$

$$\frac{P_{01}^\gamma}{P_{10}^\gamma} = \frac{n_1}{n_0} = 3 \exp(-T_*/T_\gamma) \approx 3 \left(1 - \frac{T_*}{T_\gamma}\right)$$

$$\frac{P_{01}^c}{P_{10}^c} \approx 3 \left(1 - \frac{T_*}{T_K}\right)$$

Collisions

$$\frac{P_{01}^\alpha}{P_{10}^\alpha} \approx 3 \left(1 - \frac{T_*}{T_c}\right)$$

Lyman- $\alpha$

# Equilibrium

If all three processes are active

$$n_0 [ P_{01}^\gamma + P_{01}^c + P_{01}^\alpha ] = n_1 [ P_{10}^\gamma + P_{10}^c + P_{10}^\alpha ]$$

$$\frac{n_1}{n_0} \approx 3 \left( 1 - \frac{T_*}{T_s} \right)$$

$$\frac{P_{01}^c}{P_{10}^c} \approx 3 \left( 1 - \frac{T_*}{T_K} \right)$$

$$T_s^{-1} = \frac{T_\gamma^{-1} P_{10}^\gamma + T_K^{-1} P_{10}^c + T_c^{-1} P_{10}^\alpha}{P_{10}^\gamma + P_{10}^c + P_{10}^\alpha}$$

$$T_s^{-1} = \frac{T_\gamma^{-1} + x_c T_K^{-1} + x_\alpha T_c^{-1}}{1 + x_c + x_\alpha}$$

$$x_c = \frac{P_{10}^c}{P_{10}^\gamma}$$

$$x_\alpha = \frac{P_{10}^\alpha}{P_{10}^\gamma}$$

$$P_{10}^\gamma = A_{10} + B_{10} \bar{J} \approx \frac{A_{10} T_\gamma}{T_*}$$

# Spin Temperature

$$T_S^{-1} = \frac{T_\gamma^{-1} + \tilde{x}_\alpha \tilde{T}_c^{-1} + x_c T_K^{-1}}{1 + \tilde{x}_\alpha + x_c}.$$

$$x_c = \frac{\kappa_{10} n_H T_\star}{A_{10} T_\gamma}$$

$k_{10}$

Tabulated in Zygelman 2005

$$k_{10} \approx \begin{cases} 1.19 \times 10^{-10} T_2^{0.74 - 0.20 \ln T_2} \text{ cm}^3 \text{ s}^{-1} & (20 \text{ K} < T < 300 \text{ K}) \\ 2.24 \times 10^{-10} T_2^{0.207} e^{-0.876/T_2} \text{ cm}^3 \text{ s}^{-1} & (300 \text{ K} < T < 10^3 \text{ K}) \end{cases}$$

Physics of Interstellar and Intergalactic Medium -Draine, 2011

$$\tilde{x}_\alpha = \frac{8\pi\lambda_{\text{Ly}\alpha}^2 \gamma T_\star}{9A_{10} T_\gamma} \tilde{S}_\alpha J_\alpha = 1.81 \times 10^{11} (1+z)^{-1} \tilde{S}_\alpha J_\alpha$$

Hirata 2005

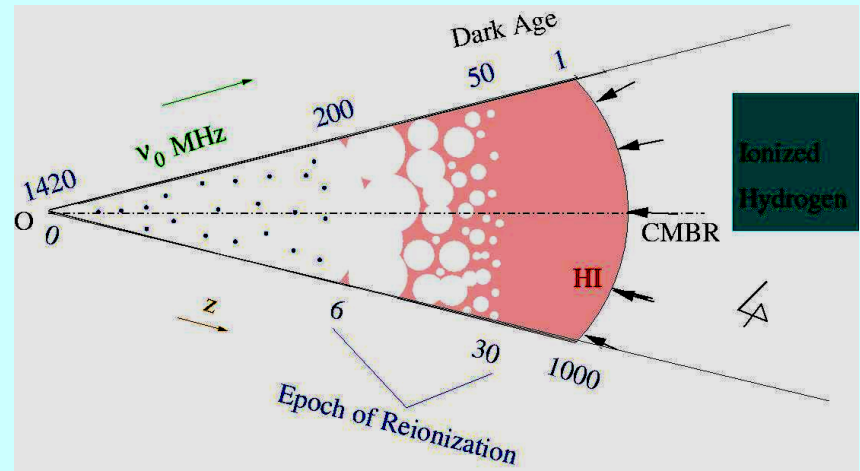
$$\gamma = 50 \text{ MHz}$$

$\tilde{S}_\alpha$  is a factor of order unity

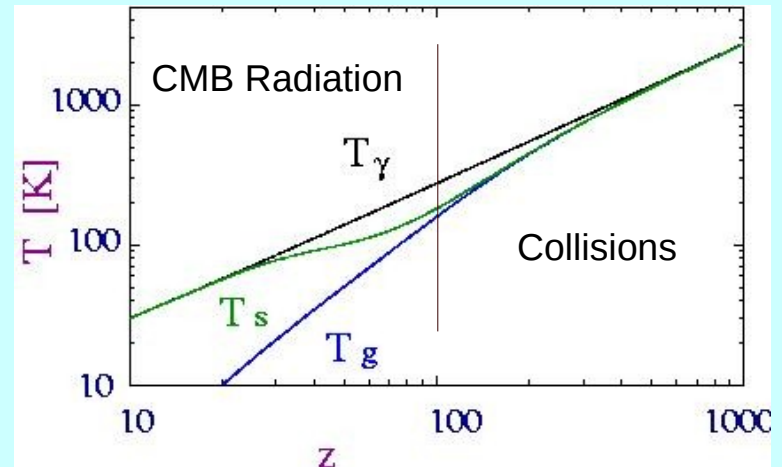
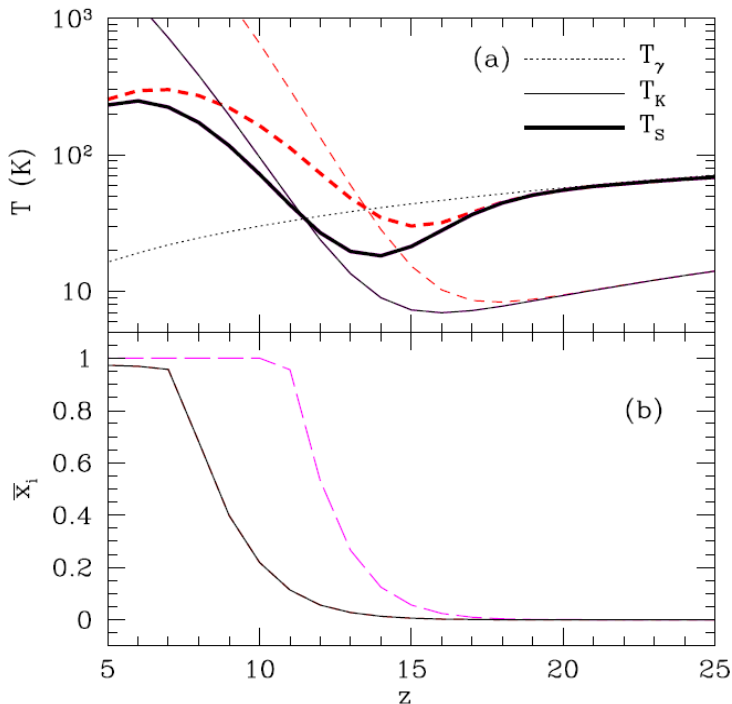
$J_\alpha$  is the flux of Ly $\alpha$  photons (in  $\text{cm}^{-2} \text{ s}^{-1} \text{ Hz}^{-1} \text{ sr}^{-1}$ ),

# Evolution of Spin Temperature

Cosmic Dawn and Epoch of Reionization  
 First stars and AGNs heat IGM  
 Emit Lyman- $\alpha$  photons - Couple  $T_s$  to  $T_K$



Dark Ages – No Luminous Sources



# Summary

- HI 21-cm line - spin-flip in ground state of neutral Hydrogen
- Potential probe of Universe over large redshift range
- Relative population ratio – quantified by spin temperature  $T_s$
- $z > 100$  - collisions couple  $T_s$  to  $T_K$  gas kinetic temperature
- $z < 100$  –  $T_s$  coupled to  $T_\gamma$  CMBR temperature
- After first stars and AGNs formed Lyman- $\alpha$  couples  $T_s$  to  $T_K$