

# Square Kilometre Array - An overview

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### SKA1 SYSTEM BASELINE2 DESCRIPTION

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Context.....AG-BD-RE  
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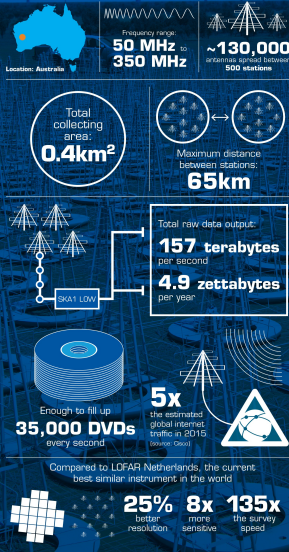


### SKA1-LOW CONFIGURATION

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Context.....PPP-PPP-PPP-TTT  
Revision.....A  
Author.....SKAO Science Team  
Date.....2015-10-28  
Document Classification.....UNRESTRICTED  
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## SKA1 LOW - the SKA's low-frequency instrument

The Square Kilometre Array (SKA) will be the world's largest radio telescope, revolutionising our understanding of the Universe. The SKA will be built in two phases - SKA1 and SKA2 - starting in 2018, with SKA1 representing a fraction of the full SKA. SKA1 will include two instruments - SKA1 MID and SKA1 LOW - observing the Universe at different frequencies.



## • Main science goals :

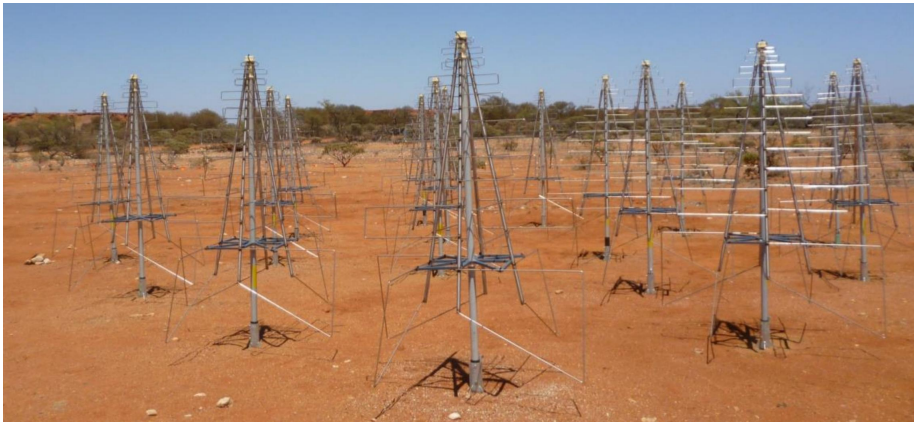
- Cosmic dawn
- Epoch of reionization
- Pulsar timing low
- Pulsar search low

## • Location : Murchison Shire of Western Australia

## • Bandwidth : 300 MHz (50 to 350 MHz frequency range)

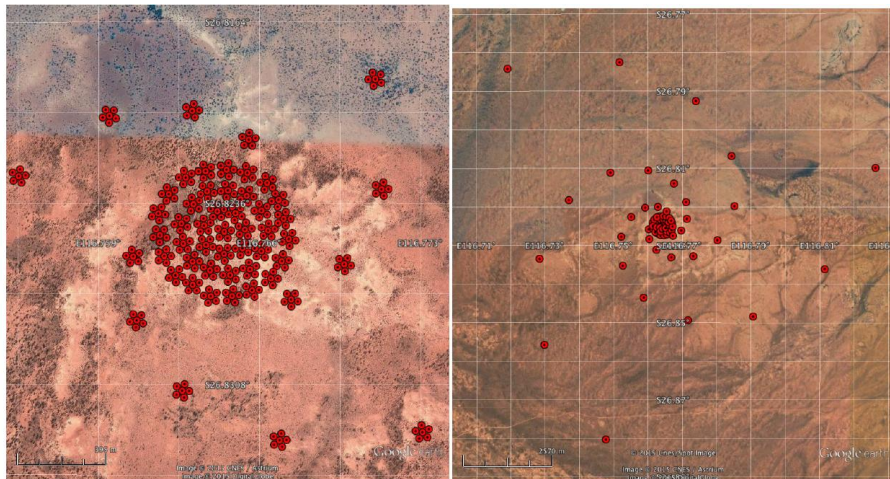
## • Will consist of an array of ~130,000 dual-polarisation antenna elements.

## • Aperture array stations (10s of metres in diameter): Antenna elements will be combined in groups so as to each act like single large antennas.



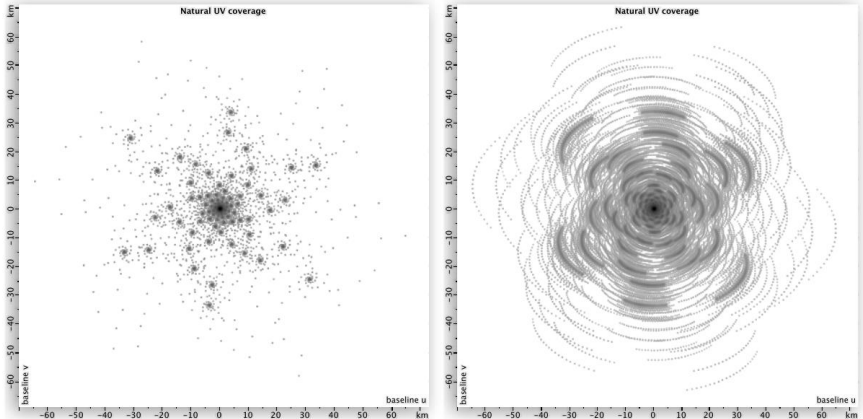
- Number of aperture array stations  $\sim 564$
- Capable of forming one or more 'beams' on the sky.
- Minimum/maximum separation between the stations:  $\sim 30$  m/65 km.
- FOV  $\sim 20$  deg<sup>2</sup>.
- Long/short dipoles at the bottom/top cover the lowest/highest frequencies

# SKA1-low antenna configuration



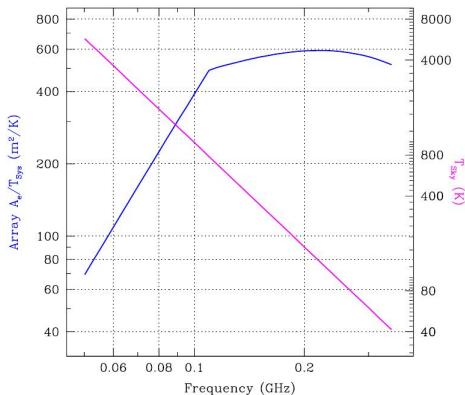
**Figure 1: The SKA1-low core (left) and skirt (right). The tightly packed system of randomised concentric rings extends to a radius of about 350m. A tightly wound three arm logarithmic spiral extends between radii of 350m – 6400m.**

# SKA1-low UV coverage



**Figure 4: The SKA1-low snap-shot (left) and 4-hour tracking (right) visibility coverage for a monochromatic observation at a nominal declination of -30.**

# Sensitivity



- System Temperature

$$T_{sys} \sim 60 \times (300 \text{ MHz}/\nu)^{2.55} \text{ K}$$

- The sensitivity curve drops steeply below 110 MHz because the sky noise, which dominates the system temperature
- Above 110 MHz the increase of collecting area as  $\lambda^2$  almost cancels the sky noise, and the ratio ( $A_e/T_{sys}$ ) is almost flat.

# How will SKA1 be better than today's best radio telescopes?



Astronomers assess a telescope's performance by looking at three factors - **resolution**, **sensitivity**, and **survey speed**. With its sheer size and large number of antennas, the SKA will provide a giant leap in all three compared to existing radio telescopes, enabling it to revolutionise our understanding of the Universe.

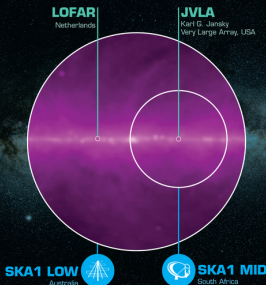


**SKA1 LOW x1.2** LOFAR NL

**SKA1 MID x4** JMLA

## RESOLUTION

Thanks to its size, the SKA will see smaller details, making radio images less blurry, like reading glasses help distinguish smaller letters.



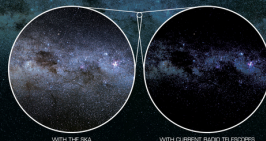
**SKA1 LOW x135** LOFAR NL

**SKA1 MID x60** JMLA

## SURVEY SPEED

Thanks to its sensitivity and ability to see a larger area of the sky at once, the SKA will be able to observe more of the sky in a given time and so map the sky faster.

The **Square Kilometre Array** (SKA) will be the world's largest radio telescope. It will be built in two phases - SKA1 and SKA2 - starting in 2018, with SKA1 representing a fraction of the full SKA. SKA1 will include two instruments - **SKA1 MID** and **SKA1 LOW** - observing the Universe at different frequencies.



**SKA1 LOW x8** LOFAR NL

**SKA1 MID x5** JMLA

## SENSITIVITY

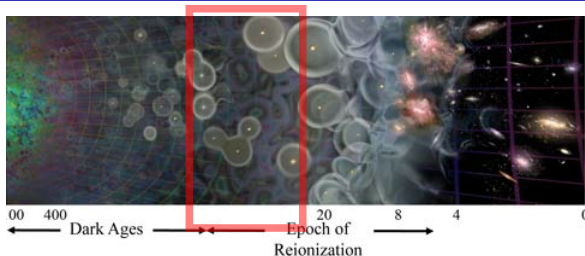
Thanks to its many antennas, the SKA will see fainter details, like a long-exposure photograph at night reveals details the eye can't see.



# Imaging the first sources during the cosmic dawn using SKA

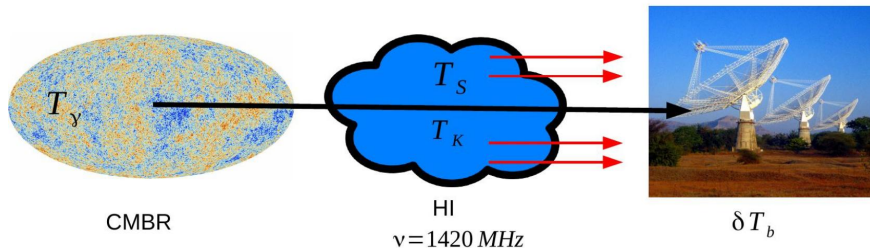
Ghara R., Choudhury T. R., Datta K. K., Choudhuri S., arXiv:1607.02779

# Motivation



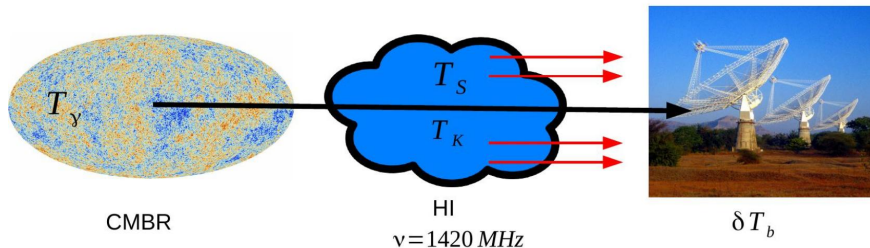
- Can we detect the first sources with radio interferometers like the SKA1-low? What informations can be obtained from that?
- IGM properties:
  - Neutral fraction
  - Kinetic temperature
- Source properties:
  - Mass of the source
  - Age
  - Escape fraction of UV photons  $f_{\text{esc}}$
  - UV and X-ray luminosity etc..

# Differential brightness temperature ( $\delta T_b$ )



$$\delta T_b = 27 x_{\text{HI}} (1 + \delta_B) \left( \frac{\Omega_B h^2}{0.023} \right) \left( \frac{0.15}{\Omega_m h^2} \frac{1+z}{10} \right)^{1/2} \left[ 1 - \frac{T_\gamma}{T_S} \right] \text{ mK}$$

# Differential brightness temperature ( $\delta T_b$ )



Density contrast

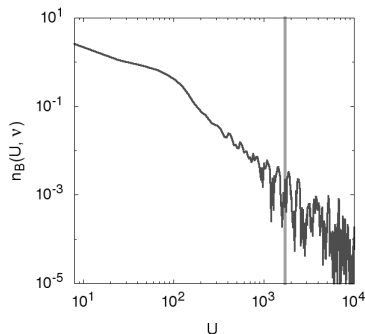
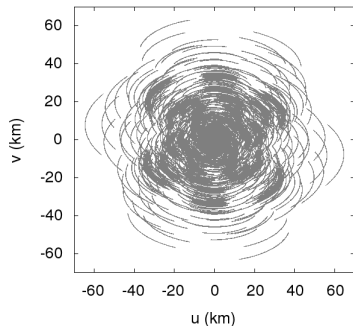
CMBR temperature

$$\delta T_b = 27 x_{\text{HI}} (1 + \delta_B) \left( \frac{\Omega_B h^2}{0.023} \right) \left( \frac{0.15}{\Omega_m h^2} \frac{1+z}{10} \right)^{1/2} \left[ 1 - \frac{T_\gamma}{T_S} \right] \text{mK}$$

neutral fraction of hydrogen

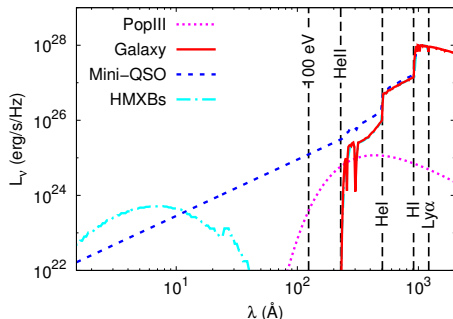
Spin temperature

# SKA1-low Baseline distribution



- Baseline  $U = d/\lambda$ .  $d$  is the distance between the antenna pair and  $\lambda$  observational wavelength.
- $n_B(U, \nu)$  is the number of antenna pairs having same baseline  $U$  at frequency  $\nu$ .

Possible sources: **popIII stars**, Galaxies, **mini-QSO**, HMXBs..

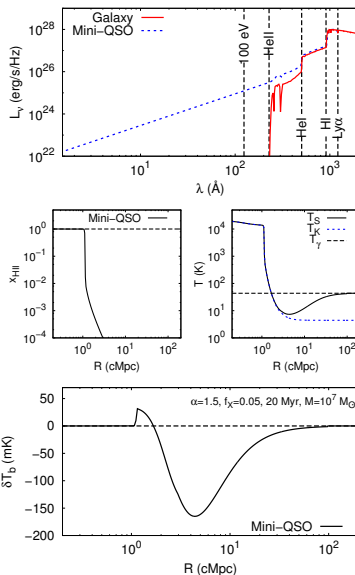


- $M_{\star} = 10^3 M_{\odot}$  for popIII  
 $= 10^7 M_{\odot}$  for others.
- Spectral index  $\alpha = 1.5$
- Ratio of X-ray and UV luminosity  $f_X = 0.05$
- $f_{\text{esc}} = 0.1$
- $t_{\text{age}} = 20 \text{ Myr}$

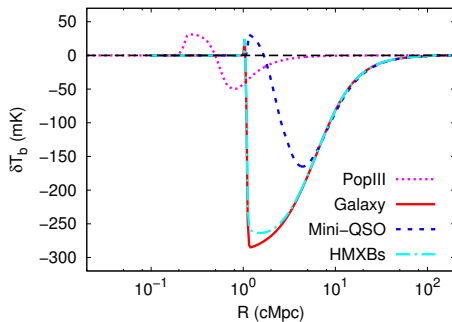
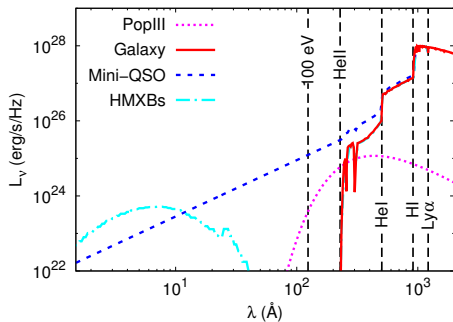
- The stellar part of the source is generated using PEGASE2 code : Salpeter IMF with 1-100  $M_{\odot}$  population II stars, metallicity  $0.001 Z_{\odot}$ .
- Mini-QSO :  $I_q \propto E^{-\alpha}$  , PopIII : black-body spectrum, HMXBs : absorption of soft-X-rays in ISM (Fragos et al 2013)

# One-dimensional radiative transfer

- Inputs
  - SED :  $L_{\gamma\alpha}$ , UV, X-rays
  - An uniform background IGM
- Generate  $x_{\text{HII}}$  and  $T_K$  profile
  - Photo-ionization, Collisional ionization, Recombinations
  - Photoelectric heating, Compton heating, cooling due to collisional excitation, Hubble expansion etc.
- Assume  $L_{\gamma\alpha}$  photon flux reduces as  $1/R^2$  with radial distance  $R$ .
- Calculate the coupling coefficients and spin temperature ( $T_S$ ) profile.
- Use  $x_{\text{HII}}$  and  $T_S$  profiles to generate  $\delta T_b$  profile.



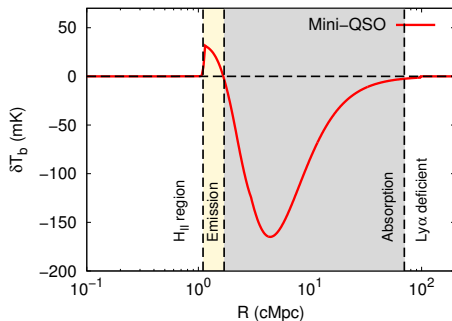
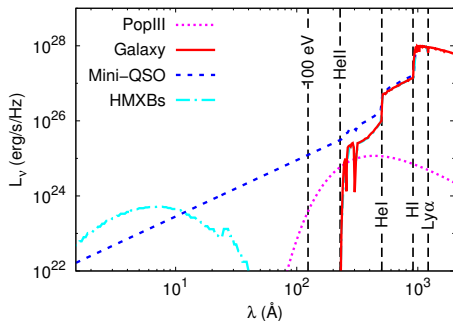
# $\delta T_b$ profiles



- Signals are different around different type of sources. Redshift = 15

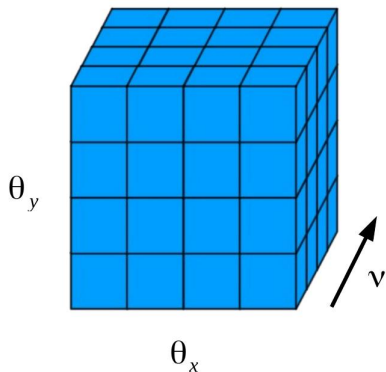


# $\delta T_b$ profiles



- Signals are different around different type of sources. Redshift = 15
- $\delta T_b$  pattern around a source :
  - H II region at the center ( $\delta T_b = 0$ ),
  - Emission region ( $\delta T_b > 0$ ),
  - Absorption region ( $\delta T_b < 0$ ),
  - Ly $\alpha$  deficient region ( $\delta T_b \sim 0$ ).

# Isolated source : simple Model



- The angular extent of the simulation box is determined by the minimum baseline.
- The angular resolution is determined by the maximum baseline.
- Length of the box along the frequency direction is determined by the observational band width.
- We assume that the source is at the center of the simulation box.
- We will image the central frequency channel.

- Differential brightness temperature :

$$\delta T_b(\vec{\theta}, \nu) \propto x_{\text{HI}}(\mathbf{x}, z)[1 + \delta_B(\mathbf{x}, z)] \left[ 1 - \frac{T_\gamma(z)}{T_S(\mathbf{x}, z)} \right]$$

- Sky specific intensity :

$$I_\nu(\vec{\theta}) = \frac{2k_B\nu^2}{c^2} \delta T_b(\vec{\theta}, \nu)$$

- Flux per synthesised beam :

$$S_\nu = I_\nu(\vec{\theta}) \times \Delta\Omega,$$

where  $\Delta\Omega = (\Delta\theta)^2$  is the beam solid angle.

- Dirty image : We first obtain the visibilities of the signal at each  $uv$  grid point and then multiply the signal with the baseline ( $uv$ ) sampling function.

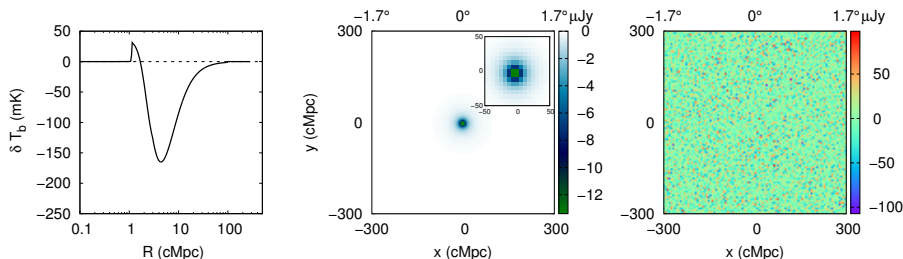
# Noise simulation

- The system noise at different baselines and frequency channels are expected to be Gaussian random variables with zero mean.
- rms noise :

$$\sqrt{\langle N^2 \rangle} = \frac{\sqrt{2} k_B T_{\text{sys}}}{A_{\text{eff}} \sqrt{\Delta \nu_c \Delta t_c}} \quad (1)$$

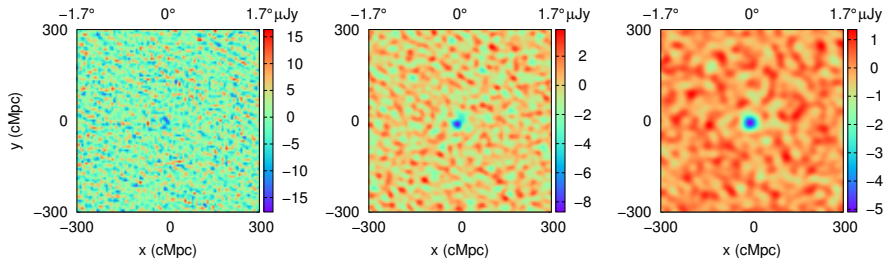
- Pixel noise can be reduced using long observation time.
- Also by a factor  $1/\sqrt{n_{\text{Base}}^{i,j}}$ , where  $n_{\text{Base}}^{i,j}$  is the number of antenna pairs associated with the  $i,j$  th pixel.
- Reduction factor :  $\sqrt{t_{\text{obs}}/t_{\text{obs}}^{uv}}$ ,  $t_{\text{obs}}^{uv} = 4 \text{ h}$
- Reverse Fourier transform to get noise maps in real space.

# Simulated signal and noise



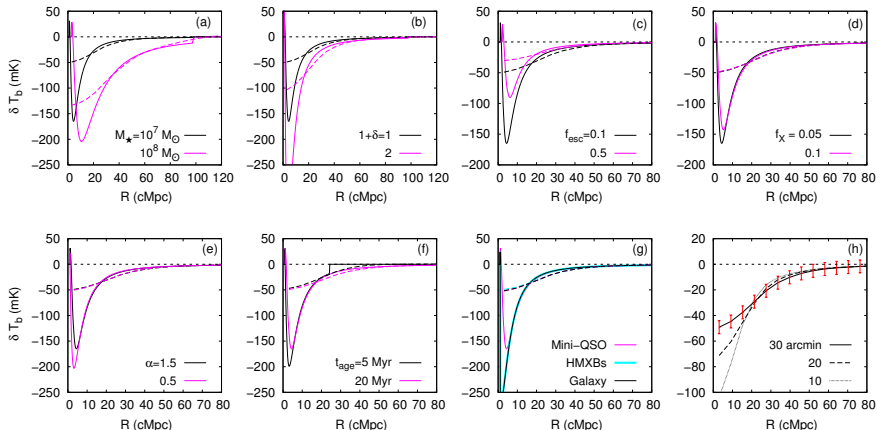
- Fiducial source :  $M_{\star} = 10^7 M_{\odot}$ , Spectral index  $\alpha = 1.5$ , ratio of X-ray and UV luminosity  $f_X = 0.05$ ,  $f_{\text{esc}} = 0.1$ ,  $t_{\text{age}} = 20 \text{ Myr}$ ,
- Angular resolution :  $2'$
- Noise map correspond to an observation time of 2000 h and frequency resolution of 100 kHz.

# Smoothed maps (Foregrounds free)



- Angular resolution :  $2'$
- Observational time : 2000 h, frequency resolution : 100 kHz
- Gaussian filter of size  $10'$  (left),  $20'$  (middle),  $30'$  (right).
- SNRs : 4 (left), 7.5 (middle), 11 (right)

# Parameter estimation?



- Fiducial source : Mini-QSO ( $M_\star = 10^7 M_\odot$ ,  $\alpha = 1.5$ ,  $f_X = 0.05$ ,  $f_{\text{esc}} = 0.1$ ,  $t_{\text{age}} = 20 \text{ Myr}$ )
- Angular resolution :  $2'$ , Gaussian filter :  $30'$

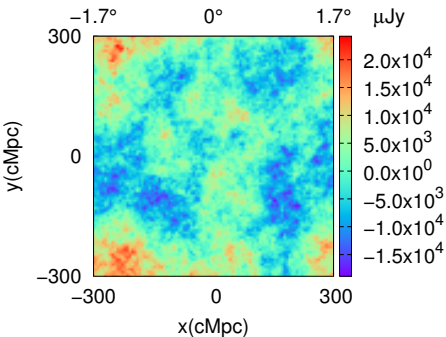
# Signal to noise ratio

Source	$M_{\star}$	$1 + \delta$	$f_{\text{esc}}$	$1 + z$	Filter	SNR1	SNR2
Mini-QSO	$10^7 M_{\odot}$	1	0.1	16	$30'$	11.1	9.1
Mini-QSO	$10^6 M_{\odot}$	1	0.1	16	$30'$	3.6	3.4
Mini-QSO	$10^8 M_{\odot}$	1	0.1	16	$30'$	25.9	20.2
Mini-QSO	$10^7 M_{\odot}$	2	0.1	16	$30'$	20.4	17.5
Mini-QSO	$10^7 M_{\odot}$	1	0.5	16	$30'$	5.2	4
Mini-QSO	$10^7 M_{\odot}$	1	0.1	11	$30'$	46	38
Mini-QSO	$10^7 M_{\odot}$	1	0.1	16	$10'$	4.2	4.0
Galaxy	$10^7 M_{\odot}$	1	0.1	16	$30'$	11.3	9.4
HMXBs	$10^7 M_{\odot}$	1	0.1	16	$30'$	11.2	9.2

- $t_{\text{obs}} = 2000 \text{ h}$ ,  $\Delta\nu = 100 \text{ kHz}$
- SNR1 : without foregrounds
- SNR2 : with foregrounds

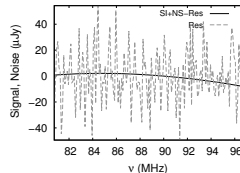
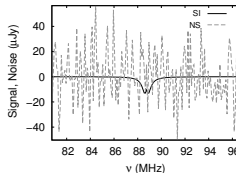
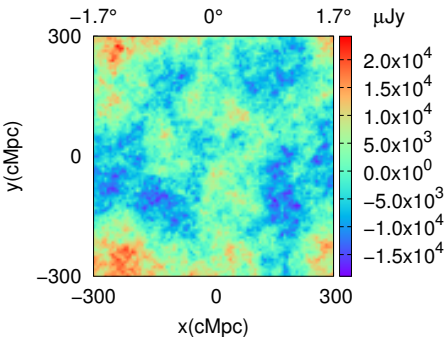


# Foregrounds?

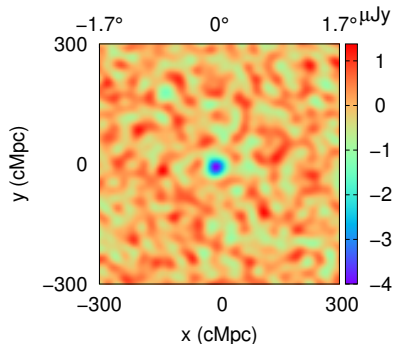


- Galactic Synchrotron radiation
- Unresolved extragalactic point sources

# Foregrounds?



- Galactic Synchrotron radiation
- Unresolved extragalactic point sources
- $\text{SNR} \sim 9$



# Using a filter

- Signal estimator:

$$\langle \hat{E} \rangle = \int d^2 U \int d\nu S(\vec{U}, \nu) S_f^*(\vec{U}, \nu) n_B(\vec{U}, \nu)$$

- Noise :

$$\langle (\Delta \hat{E})^2 \rangle_{\text{NS}} = A_{\text{NS}} \int d^2 U \int d\nu |S_f(\vec{U}, \nu)|^2 n_B(\vec{U}, \nu)$$

where,

$$A_{\text{NS}} = \langle \hat{N}^2 \rangle \frac{2}{N_{\text{ant}}(N_{\text{ant}} - 1)} \frac{\Delta \nu_c}{B_\nu}$$

- Foregrounds : (Datta et al 2007)

$$\begin{aligned} \langle (\Delta \hat{E})^2 \rangle_{\text{FG}} &= \int d^2 U \int d\nu_1 \int d\nu_2 \left( \frac{2K_B}{c^2} \right)^2 (\nu_1 \nu_2)^2 \\ &\quad \times n_B(\vec{U}, \nu_1) n_B(\vec{U}, \nu_2) C_{2\pi U}(\nu_1, \nu_2) \\ &\quad \times S_f^*(\vec{U}, \nu_1) S_f(\vec{U}, \nu_2) \end{aligned}$$

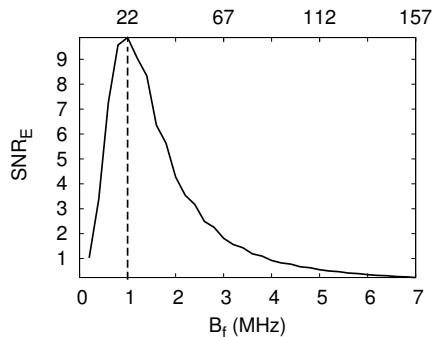
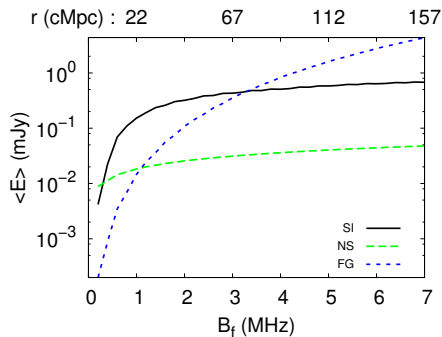
- Filter :

$$S_f(\vec{U}, \nu) = \left(\frac{\nu}{\nu_c}\right)^2 \left[ S_T(\vec{U}, \nu, B_f) - \frac{\Theta(1 - |\nu - \nu_c|/B')}{B'} \right. \\ \left. \times \int_{\nu_c - B'/2}^{\nu_c + B'/2} S_T(\vec{U}, \nu', B_f) d\nu' \right]$$

with  $B' = 2B_f$  if  $B'$  is less than  $B_\nu$ , else  $B' = B_\nu$ .

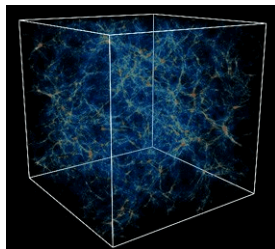
$$S_T(\vec{U}, \nu, B_f) = 0 \text{ if } |\nu - \nu_c| > \frac{B_f}{2} \\ = -1 \text{ if } |\nu - \nu_c| \leq \frac{B_f}{2}.$$

# Using a filter



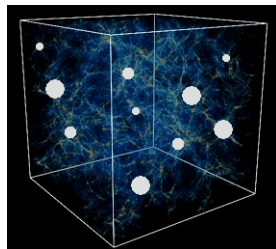
# Realistic scenario : Semi-numerical model

- Dark matter Simulation:
  - Cosmological  $N$ -body code :  
CUBEP<sup>3</sup>M ( Harnois-Deraps et al.  
2012)



# Realistic scenario : Semi-numerical model

- Dark matter Simulation:
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CUBEP<sup>3</sup>M ( Harnois-Deraps et al. 2012)
- Identify Dark matter haloes.

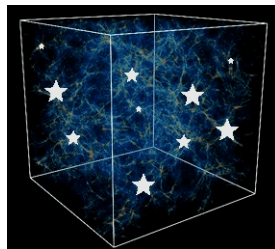


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CUBEP<sup>3</sup>M ( Harnois-Deraps et al.  
2012)
- Identify Dark matter haloes.
- These haloes are embedded with the sources of radiation.

$$M_{\star} = f_{\star} \left( \frac{\Omega_B}{\Omega_m} \right) M_{\text{halo}}$$

$f_{\star}$  is the stellar fraction of the baryon in the source.





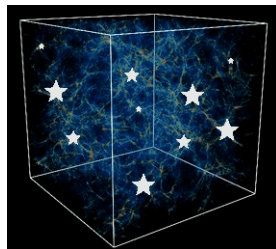
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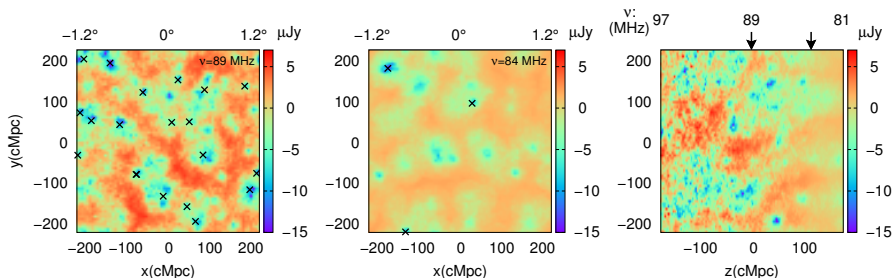
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$f_{\star}$  is the stellar fraction of the  
baryon in the source.

- Radial profile of ionization fraction  
and temperature is derived using a  
1D radiative transfer code.

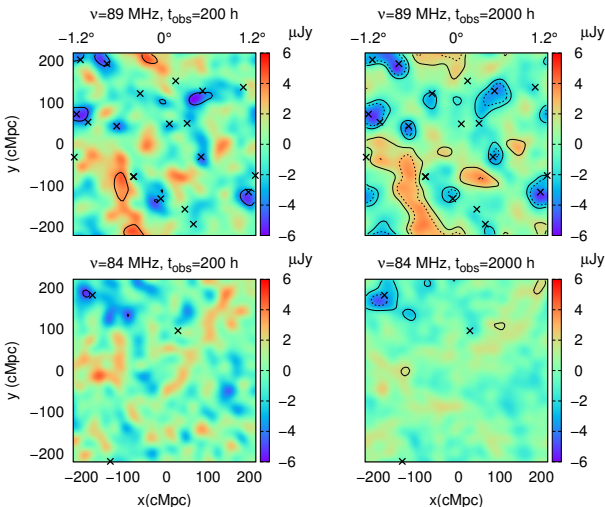


# Realistic simulation : Signal maps



- 'x' marks:  $\theta_x, \theta_y$  positions of the sources.
- Simulation :
  - simulation box of size  $300 h^{-1}$  cMpc
  - $2592^3$  particles
  - The minimum halo identified using spherical overdensity method is  $\sim 4 \times 10^9 M_\odot$ .
- Mean is subtracted at each frequency channel.

# Realistic maps : Foregrounds subtracted and smoothed



- SNR : Upper panels (4.8, 14.2), Lower panels (3.3, 10)
- **One possible observational strategy:** observe multiple fields, detect the signal in low resolution images for short  $t_{\text{obs}}$ , longer observation on the field where signal is detected, recover the  $\delta T_b$  profiles, parameter estimation.

# Conclusions

- After subtracting the foreground sufficiently and suppressing the rms noise by smoothing using a Gaussian filter, the sources are detectable with  $\sim 9 - \sigma$  confidence level over 2000 hours of observation with the SKA1-low in images of resolution  $2'$ .
- Although the 21-cm profiles around the sources get altered because of the Gaussian smoothing, the images can still be used to extract some of the source properties.
- Possible to suppress the foregrounds contributions and detect the signal using suitable filters.
- These studies are useful for making observational strategies.